

EXTENDED NEOCLASSICAL PLASMA ROTATION AND TRANSPORT THEORY

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Extended Plasma Rotation Theory(ERT)?

BACKGROUND

- Plasma rotation theories developed by two main approaches
 - Hirshman-Sigmar approach
 - Hirshman and Sigmar, Nuclear Fusion 21 (1981)
 - Most recent publication: Houlberg et al., 1998
 - most famous with two (parallel/perpendicular) Momentum Balance Equations(MBE) to calculate neoclassical rotations of multi-ions and E_r
 - Stacey-Sigmar approach
 - Stacey and Sigmar, Phys. Fluids 28, 2800(1985)
 - Most recent publication: Bae et al., Nuclear Fusion 2013
 - introduced as “Extended Plasma Rotation Theory”
 - Decomposes MBE in three coordinates (radial, poloidal, toroidal)
 - direct comparisons with V_t and V_p measurements possible
 - Radial transport calculations in radial coordinates
 - Others

BACKGROUND

■ Neoclassical plasma rotation codes

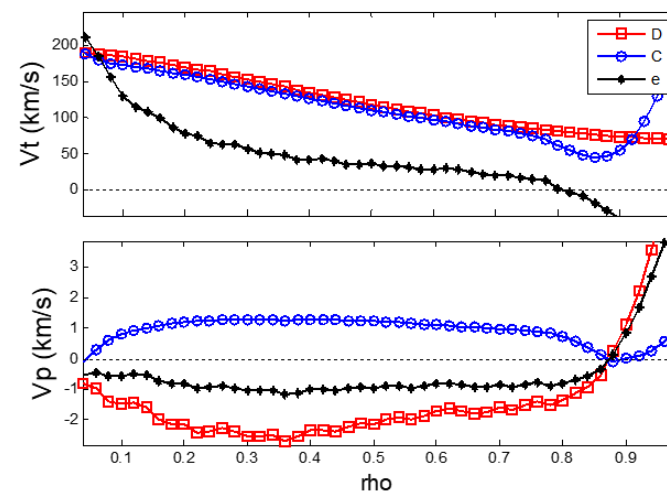
■ NCLASS

- based on Hirshman-Sigmar approach
 - Publication: Houlberg et al., Phys. Plasmas, 4 (1997)
- calculate neoclassical rotations of multi-ions
- Embedded in TRANSP

■ GTROTA

- based on Stacey-Sigmar approach
 - Publication: Bae et al., Comp. Phys. Comm. (2013)
- Uses D-shaped Miller flux surface geometry
 - Miller et. al., Phys. of Plasmas, 5 (1998)
- calculates rotation velocities up to four ion species and electron
- A non-linear iteration code in Matlab

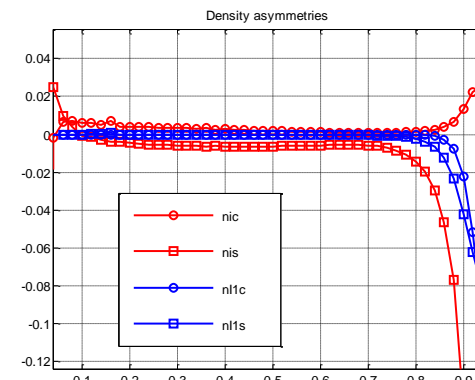
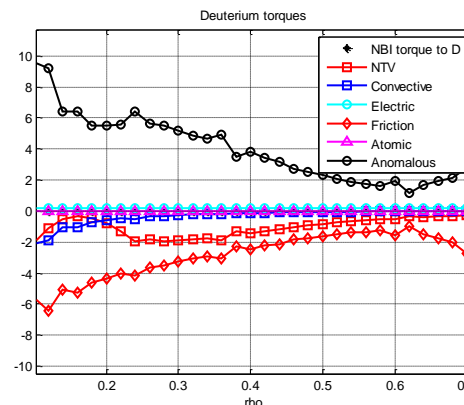
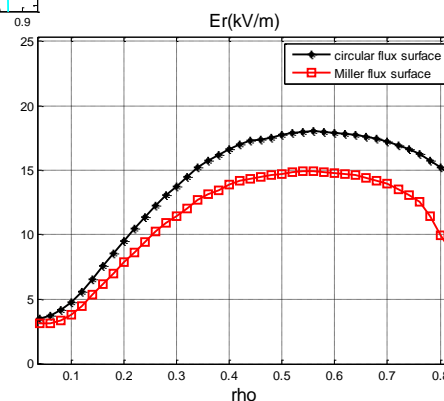
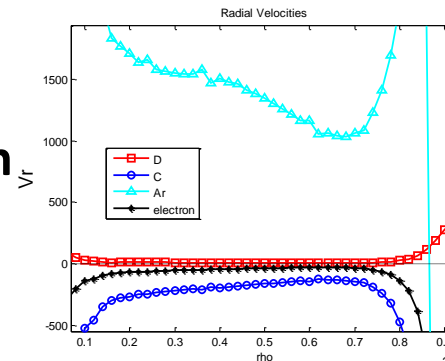
■ Others



BACKGROUND

■ GTROTA major outputs

- V_t and V_p rotation velocities of 4 ion species and electron (previous slide)
- Radial velocity (V_r) and electric field (E_r)
- All the torques in toroidal angular torque balance
 - Including viscous torques
- Poloidal in-out / up-down asymmetries
 - in density, velocity, and electrostatic potential
- V_t calculations with DIII-D, KSTAR, and EAST shots agree within 15% of carbon or Ar measurements (CES and XICS measurements) for $\rho < 0.85$



EXTENDED ROTATION THEORY

■ Based on collisionality-extended Braginskii's closure

- Meaning that it is based on Braginskii's closure but extended for arbitrary collisionality
 - Details in the next slide
- Extension to Mikhailovski-Tsyin's closure in progress [Plasma Phys 13, 785 (1971)]
 - For better accuracy in the edge Rotation Study (rho 0.85 to 1.0)

$$-\Omega_a \hat{b} \times \vec{V}_a = \partial_t \vec{V}_a + \vec{V}_a \cdot \nabla \vec{V}_a + \frac{1}{m_a n_a} \nabla \cdot \vec{\pi}_a + \frac{1}{m_a n_a} \nabla p_a - \frac{q_a}{m_a} \vec{E} - \frac{1}{m_a n_a} \vec{C}_a^{10}$$

$$-\Omega_a \left(\vec{\pi}_a \times \hat{b} - \hat{b} \times \vec{\pi}_a \right) \approx \underbrace{2 p_a \nabla \vec{V}_a}_{\text{Braginskii}} + \underbrace{\frac{4}{5} \nabla \vec{q}_a}_{\text{Mikhailovski-Tsyin}} + \text{higher order terms}$$

$$-\Omega_a \left(\hat{b} \times \vec{q}_a \right) = \frac{5}{2 m_a} p_a \nabla T_a + \frac{1}{2} \vec{C}_a^{11} + O(\text{higher})$$

EXTENDED ROTATION THEORY

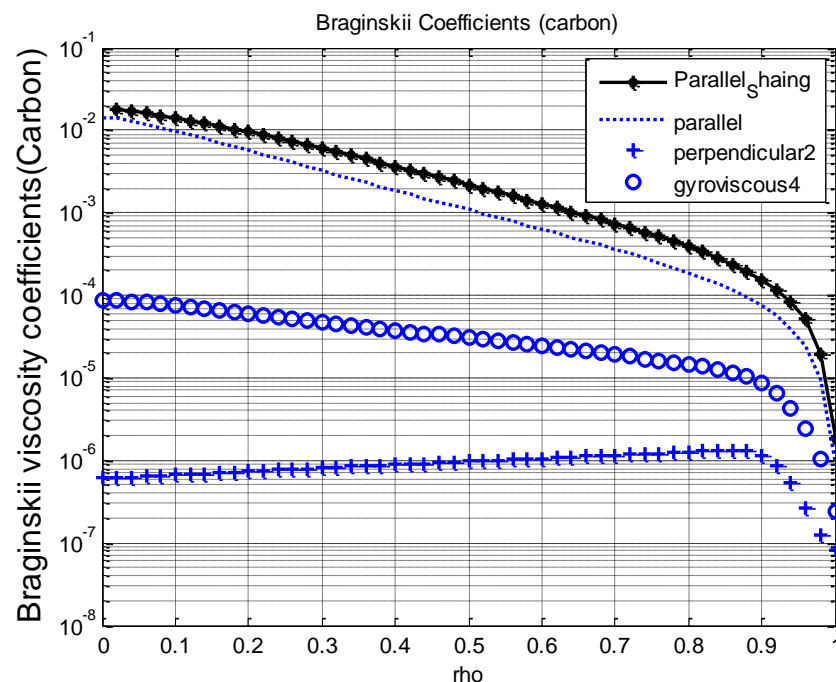
- Collisionality-extended Braginskii's viscosity representation
 - Parallel viscosity coefficient extended to low collisionality (trapped particle effect) by Shaing

-> Calculations valid for arbitrary collisionality

$$\underbrace{\eta_0 = 0.96 n T \tau}_{\text{collisionality dependent}}, \quad \underbrace{\eta_3 = \frac{1}{2} \frac{n T}{\Omega}, \quad \eta_4 = 2 \eta_3}_{\text{independent of collisionality}}, \quad \underbrace{\eta_1 = \frac{3}{10} \frac{n T}{\Omega^2 \tau}, \quad \eta_2 = 4 \eta_1}_{\text{collisionality dependent (but small contribution)}} \Rightarrow \left(\eta_0 \gg \eta_{3,4} \gg \eta_{1,2} \right)$$

- Shaing-banana plateau-PS:

$$\eta_{0j} = \frac{n_j m_j V_{thj} q R_0 \varepsilon^{-3/2} v_{jj}^*}{(1 + \varepsilon^{-3/2} v_{jj}^*)(1 + v_{jj}^*)} \equiv n_j m_j V_{thj} q R f_j$$



EXTENDED ROTATION THEORY

Retains all the terms in Toroidal and Poloidal MBE

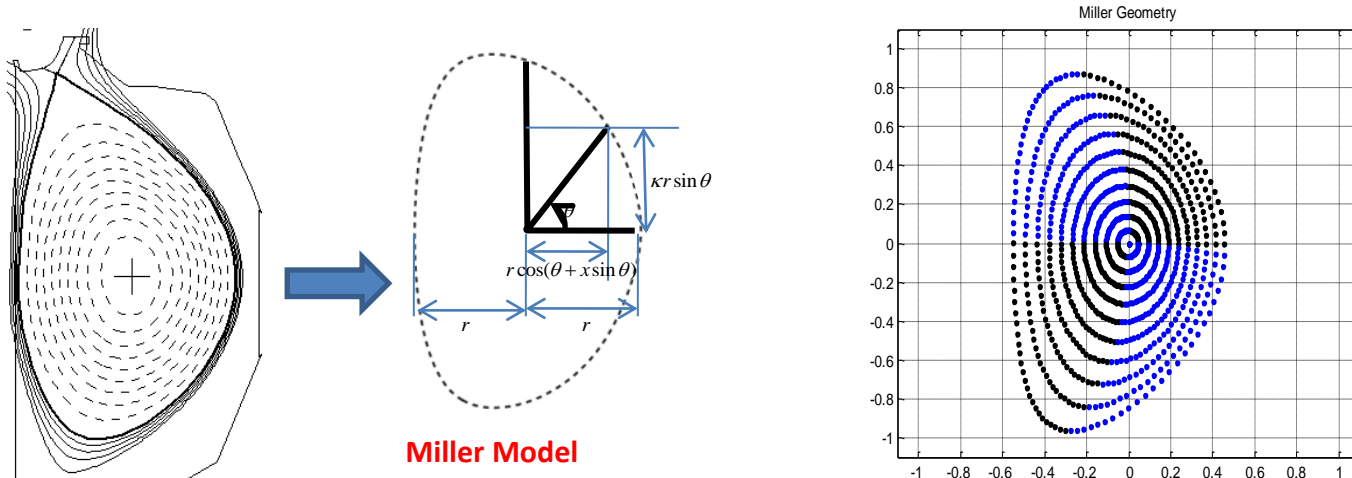
- Except vanishing ones due to equilibrium and axisymmetry (in grey)
- Includes Reynolds Stress terms (Convective and Atomic torques)
- Atomic term calculated with TRANSP
- No gyroviscous cancellation assumed
- Numerical model becomes extremely non-linear

$$\begin{aligned}
 &\text{Accelerational} \quad \boxed{\text{Convective}} \quad \text{Pressure} \quad \boxed{\text{Viscous}} \\
 &\left\langle Rm_a n_a \frac{\partial V_{a\phi}}{\partial t} \right\rangle + \left\langle Rn_a m_a (\vec{V}_a \cdot \nabla) V_{a\phi} \right\rangle + \left\langle R \frac{1}{h_\phi} \frac{\partial p_a}{\partial \phi} \right\rangle + \left\langle R (\nabla \cdot \vec{\pi}_a)_\phi \right\rangle = \\
 &\left\langle Rn_a e_a \left(-\frac{1}{h_\phi} \frac{\partial \Phi}{\partial \phi} - \frac{\partial A_\phi}{\partial t} \right) \right\rangle + \left\langle Rn_a e_a (V_{a\psi} B_p - V_{a\phi} B_\psi) \right\rangle + \left\langle RF_{a\phi}^1 \right\rangle + \left\langle RS_{a\phi}^1 \right\rangle + \underbrace{\left(-\left\langle Rm_a V_{a\phi} S_a^0 \right\rangle \right)}_{-\langle Rn_j m_j V_{atomj} V_{\phi j} \rangle} \\
 &\text{Electric} \quad \quad \quad \text{V cross B} \quad \quad \quad \text{Frictional} \quad \quad \text{External} \quad \quad \text{Atomic}
 \end{aligned}$$

$$\begin{aligned}
 &\left\langle rm_a n_a \frac{\partial V_{a\theta}}{\partial t} \right\rangle + \left\langle rn_a m_a (\vec{V}_a \cdot \nabla) V_{a\theta} \right\rangle + \left\langle r \frac{1}{h_\theta} \frac{\partial p_a}{\partial \theta} \right\rangle + \left\langle r [\nabla \cdot \vec{\pi}_a]_\theta \right\rangle = \\
 &\left\langle r \left(-\frac{1}{h_p} \frac{\partial \Phi}{\partial \theta} - \frac{\partial A_\theta}{\partial t} \right) \right\rangle + \left\langle rn_a e_a (V_{a\phi} B_r - V_{a\theta} B_\phi) \right\rangle + \left\langle r F_{a\theta}^1 \right\rangle + \left\langle r (S_{a\theta}^1 - m_a V_{a\theta} S_a^0) \right\rangle
 \end{aligned}$$

EXTENDED ROTATION THEORY

- applies D-shaped flux surface geometry using Miller model [Miller et. al., Phys. of Plasmas, 5 (1998)] with Shafranov shifts



- Flux surface averages (FSAs) and geometric scale factors(h's) for Miller geometry
 - Example: Er formula below shows FSAs and scale factors (h)

$$\bar{E}_r = -V_{thj} \bar{B}_\theta \left[V_{\theta j} \frac{\left\langle \frac{1}{1 + \varepsilon \cos \xi} \right\rangle}{\left\langle \frac{1}{h_r} \right\rangle} - V_{\phi j} \left(1 + \frac{\partial R_0(r)}{\partial r} \right) \frac{\left\langle \frac{1}{(1 + \varepsilon \cos \xi) h_r} \right\rangle}{\left\langle \frac{1}{h_r} \right\rangle} - \frac{1}{V_{thj}} \frac{1}{n_j e_j \bar{B}_\theta} \frac{\partial \bar{P}_j}{\partial r} \right]$$

EXTENDED ROTATION THEORY

- Includes first-order poloidal perturbations (poloidal asymmetries) in density, velocity, and electrostatic potential

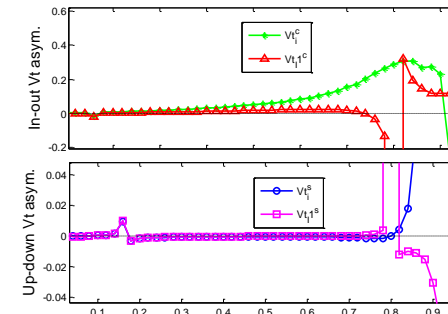
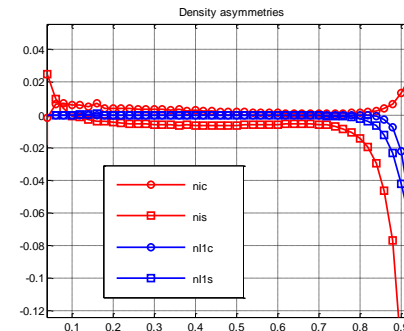
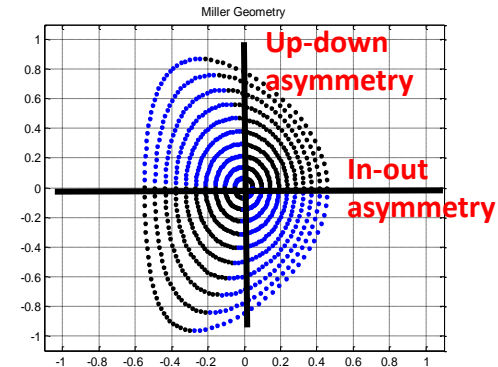
- Represented with the 1st order Fourier series

$$X_j(\rho, \phi) \approx \bar{X}_j(\rho) \left[1 + \sum_{n=1}^{\infty} \left(\underbrace{X_j^{nc}(\rho) \cos(n\phi)}_{\text{in-out asymmetry}} + \underbrace{X_j^{as}(\rho) \sin(n\phi)}_{\text{up-down asymmetry}} \right) \right]$$

- Sine function representing up-down asymmetry
- Cosine function representing in-out asymmetry

- Calculated asymmetries with Miller geometry published in two papers

- Bae et al., Nuclear Fusion 2013
- Bae et al., Phys of Plasmas, 2014
- A draft under review



- Georgia Tech Fusion Research Center has recently developed a code to calculate rotations with 10th order perturbations for accuracy
 - R. King, MS thesis, Georgia Institute of Technology, May 2019

EXTENDED ROTATION THEORY

■ Radial transport of all ion species calculated

■ Er calculated self-consistently

$$\bar{E}_r = -V_{thj} \bar{B}_\theta \left[V_{\theta j} \frac{\left\langle \frac{1}{1 + \varepsilon \cos \xi} \right\rangle}{\left\langle \frac{1}{h_r} \right\rangle} - V_{\phi j} \left(1 + \frac{\partial R_0(r)}{\partial r} \right) \frac{\left\langle \frac{1}{(1 + \varepsilon \cos \xi) h_r} \right\rangle}{\left\langle \frac{1}{h_r} \right\rangle} - \frac{1}{V_{thj}} \frac{1}{n_j e_j \bar{B}_\theta} \frac{\partial \bar{P}_j}{\partial r} \right]$$

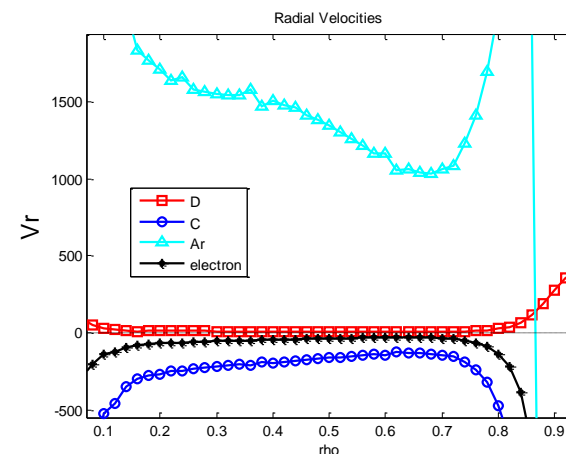
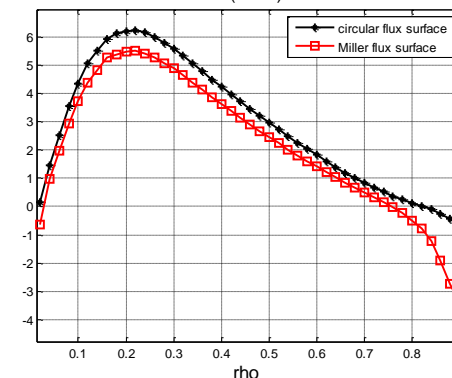
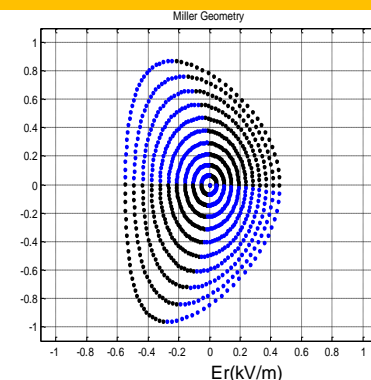
■ Calculation with Miller flux surfaces published in 2014

- Bae et al., Phys of Plasmas, 2014
- Difference in Er with Circular vs. Miller model published

$$\bar{E}_r^{cir} = \frac{1}{n_j e_j} \frac{\partial \bar{P}_j}{\partial r} - [V_{\theta j} \bar{B}_\phi - V_{\phi j} \bar{B}_\theta]$$

■ Recent investigations on Vr and Radial flux of all ion species

- calculated non-self-consistently
- GTROTA updated to investigate these calculations
- Collaboration with other researches based on radial diffusivity coefficients (later slides)



MAJOR DEVELOPMENTS

**Future directions of ERT and GTROTA
(mostly for edge rotation study)**

Advises/ideas/recommendations welcomed!

MAJOR DEVELOPMENTS

■ Extension of ERT to Mikhailovskii-Tsy-pin's closure

- Include Heat Equation
- Important for edge rotation/transport study

- Viscosity evolution equation:

$$-\Omega_a \left(\hat{b} \times \vec{\pi}_a - \vec{\pi}_a \times \hat{b} \right) = 2p_a \overline{\nabla \vec{V}_a} + \frac{4}{5} \overline{\nabla \vec{h}_a} - \vec{C}_a^{20} + \text{higher orders}$$

- Heat flow evolution equation:

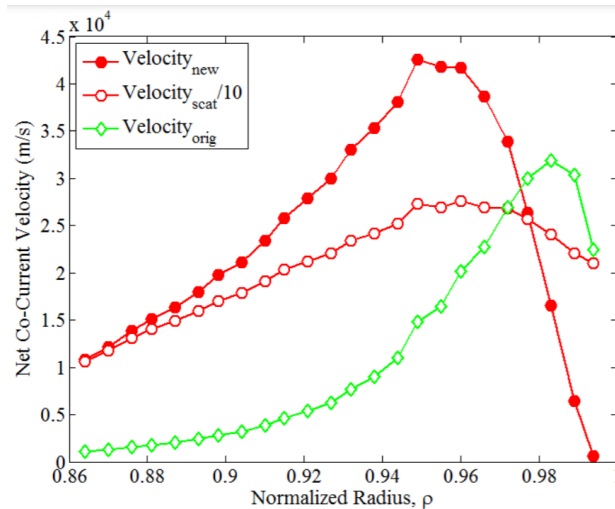
$$-\Omega_a \left(\hat{b} \times \vec{h}_a \right) = \frac{5}{2m_a} p_a \nabla T_a + \frac{1}{2} \vec{C}_a^{11} + \text{higher orders}$$

➤ This development in progress (2016 - ongoing)

- 1ST order Fourier expansion of temperature being applied
- This is next in line for the future development of ERT and GTROTA
 - doable within a few years to finalize the numerical model and implement it in GTROTA

MAJOR DEVELOPMENTS

- To add more edge physics including intrinsic rotation
 - Intrinsic co-current deuterium rotation due to Ion Orbit Loss (IOL)
 - Stacey, Phys of Plasmas 25, 122506 (2018)
 - The predominant IOL of CTR-current ions leaves a predominantly CO-current edge intrinsic rotation



- This effort shouldn't take long but will need to add other edge rotation mechanisms for higher accuracy
 - Ideas/suggestions welcomed

MAJOR DEVELOPMENTS

■ Develop an NTV theory based on Stacey-Sigmar approach

■ By considering non-axisymmetry in the formalism

■ General non-axisymmetric formalism published: Stacey and Bae, PoP (2015)

$$n_j m_j \left\langle \left[R \left(\vec{V}_j \cdot \nabla \right) \vec{V}_j \right]_\phi \right\rangle + \left\langle \left[R \left(\nabla \cdot \vec{\pi}_j \right) \right]_\phi \right\rangle = \langle R n_j e_j E_\phi^A \rangle + \langle R n_j e_j V_{rj} B_\theta \rangle + \langle R F_{\phi j} \rangle + \langle R M_{\phi j} \rangle - \langle R n_j m_j V_{atomj} V_{\phi j} \rangle$$

NTV Torque

$$\left\langle R \left(\nabla \cdot \vec{\pi} \right)_\phi \right\rangle = \left\langle R^2 \nabla \phi \cdot \left(\nabla \cdot \vec{\pi} \right) \right\rangle = \left\langle \frac{1}{R h_p} \frac{\partial \left(R^2 h_p \pi_{\psi\phi} \right)}{\partial l_\psi} \right\rangle$$

$$\eta_0 \gg \eta_{3,4} \gg \eta_{1,2}$$

|| gv ⊥

■ Dedicated presentation slides in Appendix A

$$\left\langle R \left(\nabla \cdot \vec{\pi} \right)_\phi \right\rangle = \left\langle R^2 \nabla \phi \cdot \left(\nabla \cdot \vec{\pi} \right) \right\rangle = \underbrace{\left\langle \frac{1}{R h_p} \frac{\partial \left(R^2 h_p \pi_{\psi\phi}^0 \right)}{\partial l_\psi} \right\rangle}_{\text{non-axisymmetric}} + \underbrace{\left\langle \frac{1}{R h_p} \frac{\partial \left(R^2 h_p \pi_{\psi\phi}^{34} \right)}{\partial l_\psi} \right\rangle + \left\langle \frac{1}{R h_p} \frac{\partial \left(R^2 h_p \pi_{\psi\phi}^{12} \right)}{\partial l_\psi} \right\rangle}_{\text{axisymmetric}}$$

$$\pi_{\psi\phi}^0 = -\eta_0 W_{\psi\phi}^0$$

$$W_{\psi\phi}^0 = \frac{3}{2} f_\psi f_p H^0$$

$$f_\psi \equiv \frac{B_\psi}{B}, \quad f_p \equiv \frac{B_p}{B}, \quad f_\phi \equiv \frac{B_\phi}{B}$$

MAJOR DEVELOPMENTS

■ Develop an NTV theory based on Stacey-Sigmar approach

- Next step is to develop a numerical model for GTROTA

$$\pi_{\psi\phi}^0 = -\eta_0 W_{\psi\phi}^0 \quad W_{\psi\phi}^0 = \frac{3}{2} f_{\psi} f_p H^0 \quad f_{\psi} \equiv \frac{B_{\psi}}{B}, \quad f_p \equiv \frac{B_p}{B}, \quad f_{\phi} \equiv \frac{B_{\phi}}{B}$$

■ Question is on how to best represent magnetic perturbations

- Up to the 4th order Fourier series being considered but still contemplating
- Ideas being formulated from my experimental experiences at KSTAR
 - as a team member for the ELM suppression investigation using RMP coils at KSTAR

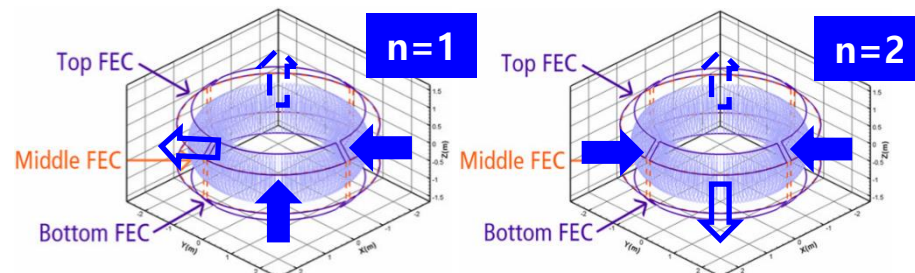
- Publication: Jayhyun Kim et al., Nucl. Fusion 57, 022001 (2017)

$$\vec{B}_{\text{total}} = \vec{B}_{\text{eq}} + \cancel{\vec{B}_{\text{int}}} + \vec{B}_{\text{RMP}} + \vec{B}_{\text{plasma response}}$$

n=2 (mid) + n=1 (top/bot)

Port	L	P	D	H
ϕ	0°	90°	180°	270°
Top	+	-	-	+
Mid	+	-	+	-
Bot	+	+	-	-

→ Proxy (n=1 NA field)
⇒ Reference (n=2 NA field)
→ Proxy (n=1 NA field)



MAJOR DEVELOPMENTS

■ Investigate the converging order of poloidal asymmetries

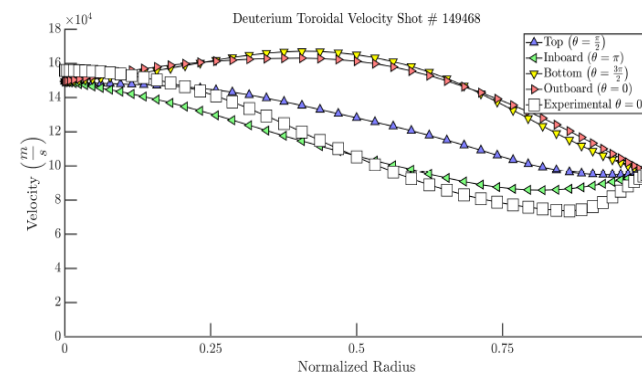
- Georgia Tech Fusion Research Center has calculated rotations with poloidal asymmetries up to 10th order

- Publication: R. King, MS thesis, Georgia Institute of Technology

$$n = n_0 + \sum \sum a_{ij} J_i(\lambda_{ij} \rho) \cos(j\theta) + b_{ij} J_i(\lambda_{ij} \rho) \sin(j\theta)$$

$$v_t = v_{t0} + \sum \sum a_{ij} J_i(\lambda_{ij} \rho) \cos(j\theta) + b_{ij} J_i(\lambda_{ij} \rho) \sin(j\theta)$$

$$\phi = \phi_0 + \sum \sum a_{ij} J_i(\lambda_{ij} \rho) \cos(j\theta) + b_{ij} J_i(\lambda_{ij} \rho) \sin(j\theta)$$



- A separate set of codes written with Mathematica and Fortran
 - Could be a verification opportunity on which order is accurate enough
 - This investigation of the appropriate converging order may take long
- Two track development approaches with GTROTA
 - Accuracy version (with higher order poloidal asymmetries)
 - Plasma control version (for near-real-time calculations of all species)
 - This will have to be collaboration with PCS developers

Ongoing Researches

- Publication plans
 - Researches under investigation/collaboration
- ❖ Please note that any research topics discussed in this talk can be considered as collaboration opportunities for others.

Publication Plans

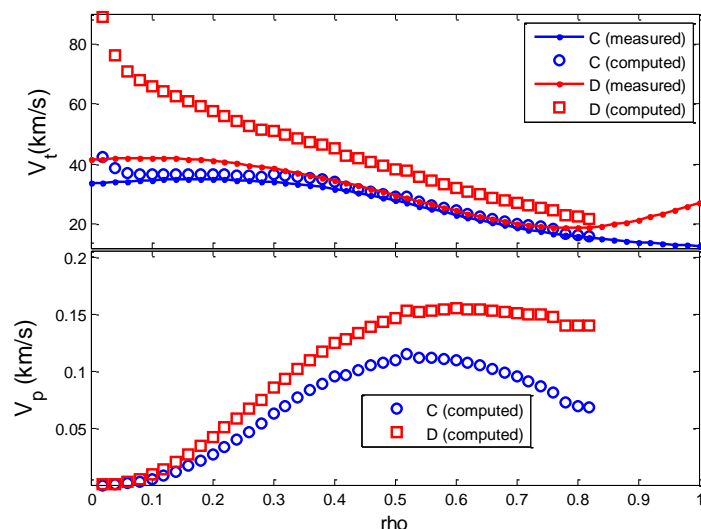
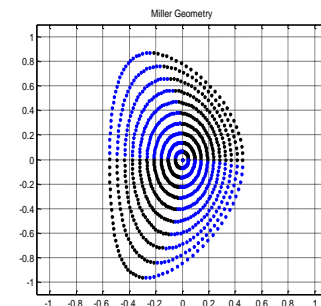
■ Theory verification with DIII-D deuterium velocity measurements in L-mode (2013-2015)

■ First test opportunity of the theory against deuterium measurements

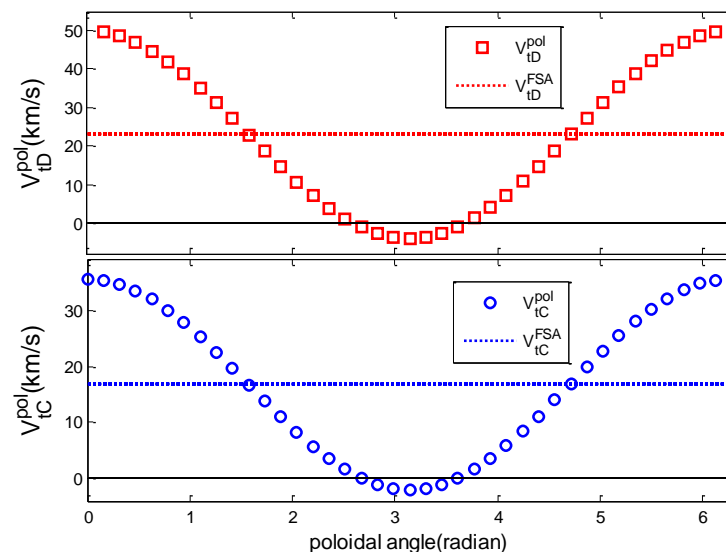
- Deuterium measurement became available from Dr. Brian Grierson
- Below: DIII-D 145180 (1220ms)

■ Interesting findings on poloidal variation of toroidal velocities

- Similar experimental findings with ECEI rotation images analyzed by G.S. Yoon at POSTECH



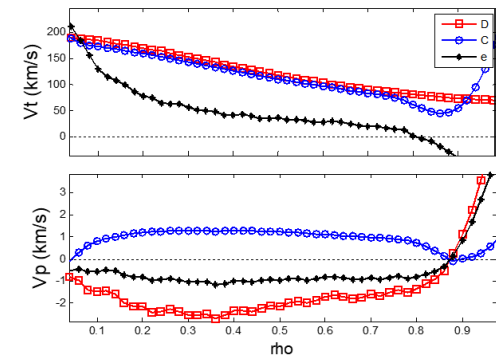
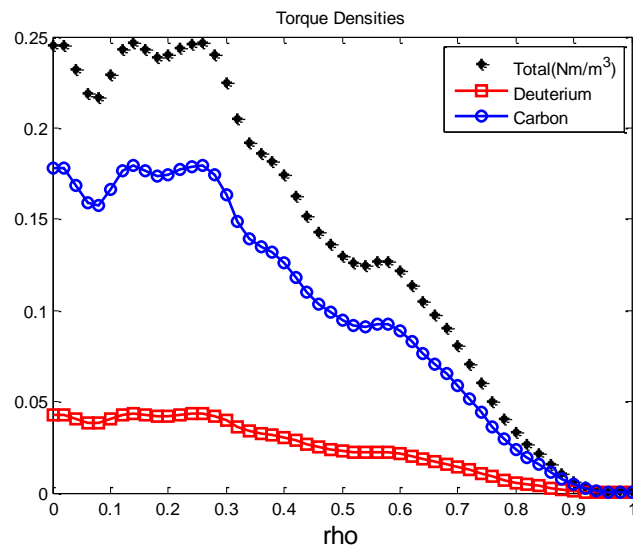
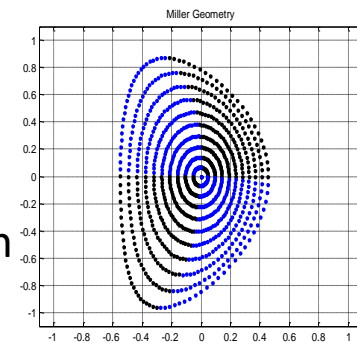
Toroidal velocities (D and C)



Toroidal velocity asymmetries (D and C) at rho=0.8

Publication Plans

- **EAST and KSTAR H- and L-mode shot analyses (2015-2017)**
 - Goal was to analyze more L-mode shots and upgrade GTROTA capability to calculate rotations with RF heatings
 - EAST shots with LHCD and ICRF shots
 - KSTAR L-mode shots
 - **Analyses finalized**
 - Need to extend GTROTA features to add HFS-LFS plotting option
 - Planning to publish with new features available in GTROTA



Publication Plans

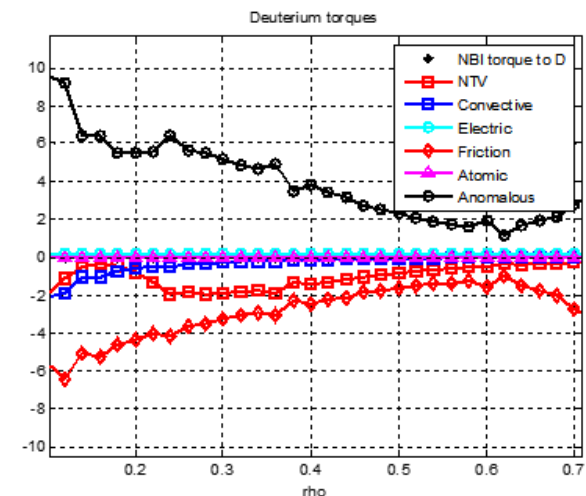
■ Theoretical investigation of Gyroviscous cancellation validity in Tokama plasmas (2014-2015)

- The well-known gyroviscous cancellation in sheared slab geometry [Plasma Physics **4**, 1766 (1992)] has been investigated
 - using a systematic perturbative method
 - based on the Mikhailovskii-Tsypin's closure relation
 - on the large gyrofrequency ordering for flowing plasmas
- $O(\delta^1)$ 1st order surviving terms (in colors):

$$\partial_t \vec{V}_{0\parallel} + \vec{V}_{d(1)} \cdot \nabla \vec{V}_{(0)\parallel} + \vec{V}_{(0)} \cdot \nabla \vec{V}_{(1)\parallel} + \frac{1}{m_a n_0} \nabla_{\parallel} p_0 = -\Omega_a \left(\hat{b} \times \vec{V}_1 \right)_{\parallel} + \frac{q_a}{m_a} \vec{E}_{0\parallel} + \frac{1}{m_a n_0} \vec{C}_{0\parallel}^{10}$$

- These are the surviving terms at one order lower than investigated by Chang and Callen [Plasma Physics **4**, 1766 (1992)]

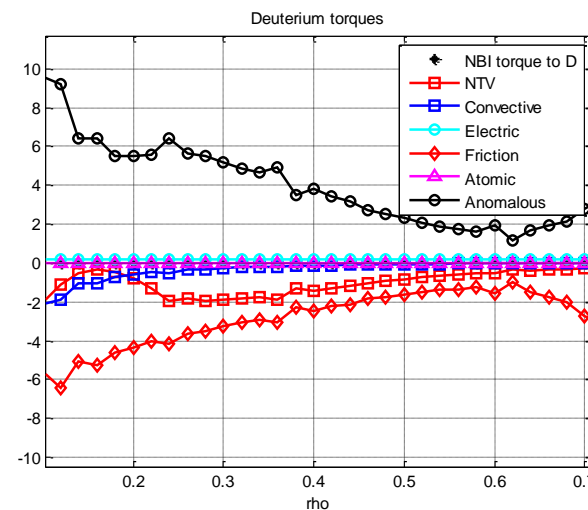
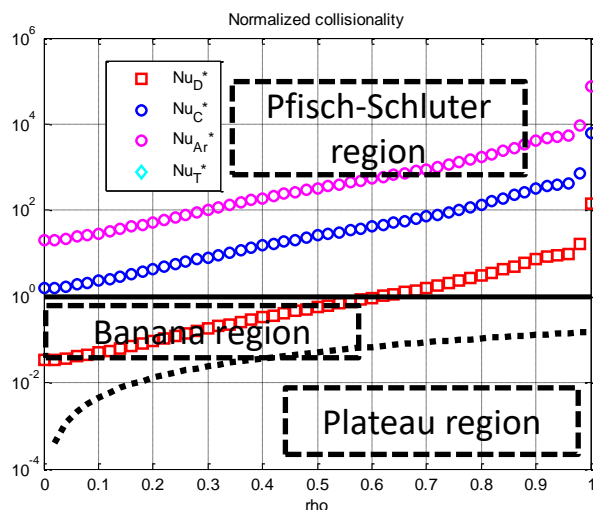
- Submitted for publication but withdrawn for numerical verifications
- Presented at conferences



Researches under Investigation/Collaboration

Collisionality effects on rotation and transport

- Motivated by too many assumptions applied in plasma researches
 - ERT retains all terms in the MBE (both toroidal and poloidal coordinates)
- This research is to holistically understand collisionality effects
 - **within a discharge**: collisionality regimes change and differ for different species
 - **among different discharges**: Different discharges have different regime distributions
- **Q: Can I investigate collisionality effects on various shots and find any consistency on rotation/transport?**
 - Will require analyses of many discharges to answer this question



Researches under Investigation/Collaboration

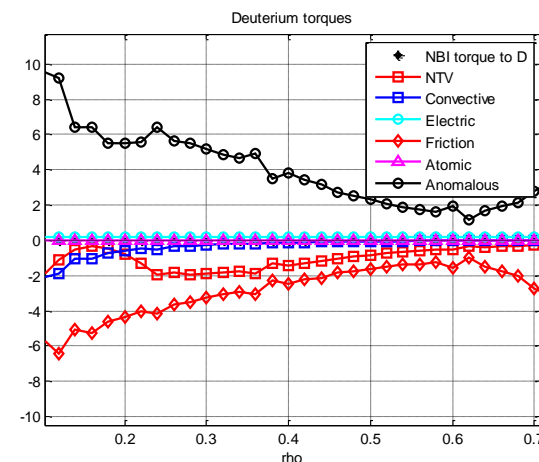
Collisionality effects on rotation and transport

- However, I can provide my rotation/transport calculations to others so that they can identify right assumptions based on the collisionality and apply appropriate theories
- Calculations in toroidal angular torques are available in GTROTA

$$\begin{aligned}
 & \left\langle R n_a n_a \frac{\partial V_{a\phi}}{\partial t} \right\rangle + \left\langle R n_a m_a (\vec{V}_a \cdot \nabla) V_{a\phi} \right\rangle + \left\langle R \frac{1}{h_\phi} \frac{\partial p_a}{\partial \phi} \right\rangle + \left\langle R (\nabla \cdot \vec{\pi}_a)_\phi \right\rangle = \\
 & \left\langle R n_a e_a \left(-\frac{1}{h_\phi} \frac{\partial \Phi}{\partial \phi} - \frac{\partial A_\phi}{\partial t} \right) \right\rangle + \left\langle R n_a e_a (V_{a\psi} B_p - V_{a\phi} B_\psi) \right\rangle + \left\langle R F_{a\phi}^1 \right\rangle + \left\langle R S_{a\phi}^1 \right\rangle + \underbrace{\left(-\left\langle R m_a V_{a\phi} S_a^0 \right\rangle \right)}_{-\langle R n_j m_j V_{atomj} V_{\phi j} \rangle}
 \end{aligned}$$

Accelerational
Convective
Pressure
Viscous

Electric
V cross B
Frictional
External
Atomic



- Plan to do the same code upgrade for poloidal torque balance in the future

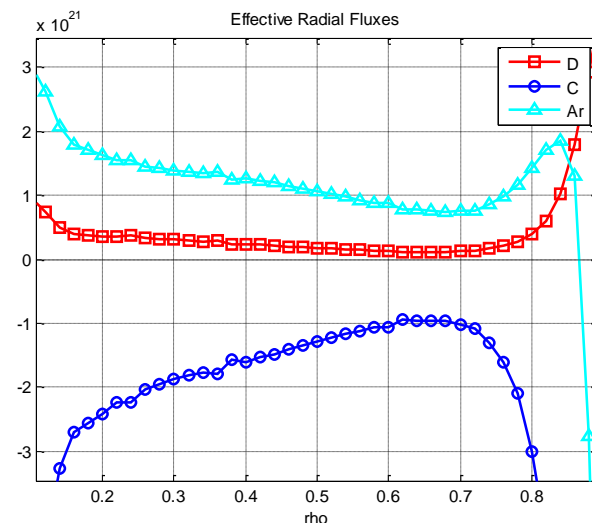
$$\left\langle r n_j m_j (\vec{V}_j \cdot \nabla) \vec{V}_{j\theta} \right\rangle + \left\langle \frac{r}{h_\theta} \frac{\partial p_j}{\partial \theta} \right\rangle + \left\langle r (\nabla \cdot \vec{\pi}_j)_\theta \right\rangle = \left\langle r F_{\theta j} \right\rangle + \left\langle r n_j e_j E_\theta \right\rangle - \left\langle r n_j e_j V_{rj} B_\phi \right\rangle + \left\langle r S_{\theta j}^1 \right\rangle - \left\langle r m_j V_{\theta j} S_j^0 \right\rangle$$

Researches under Investigation/Collaboration

- **A coupled study of ion species rotation and transport**
 - Motivation is to compare radial transport (V_r and flux) calculations from GTROTA to other calculations based on diffusivity coefficients
 - Collaboration with KAIST team on Ar transport study (2015 – 2017)
 - with density gradient: introduces “Diffusion coefficient (D)”
 - with convective effect: introduces “Convective coefficient (V)”

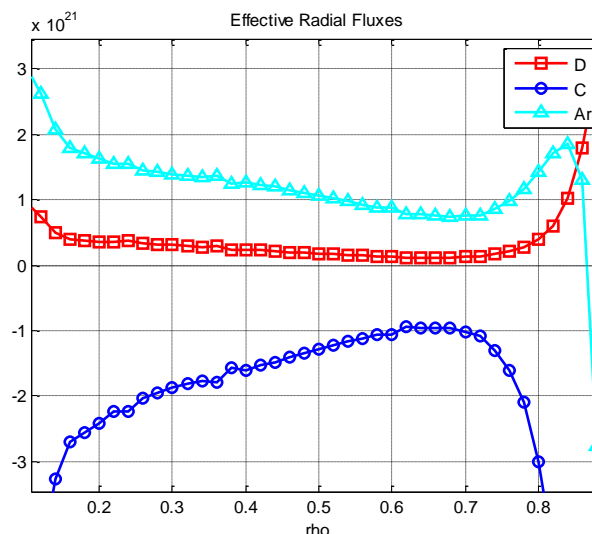
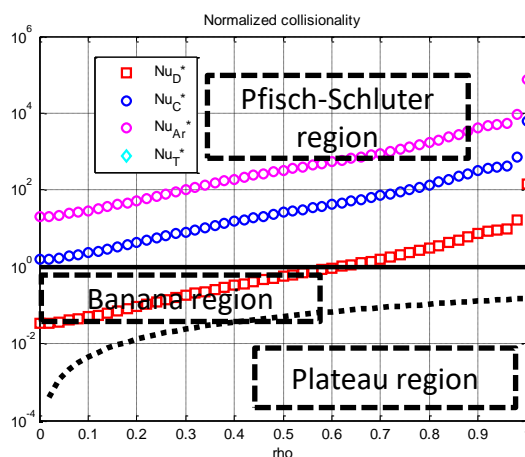
$$\frac{\partial n_a}{\partial t} + \underbrace{\nabla \cdot (n_a \vec{V}_a)}_{\text{D}fn(n_a) + \text{V}fn(\vec{V}_a)} = S_a^0$$

- KSTAR experiments with Ar injection
- Analyses delayed due to the temperature measurement accuracy
- **A dedicated conference slides in Appendix B**



Researches under Investigation/Collaboration

- A coupled study of ion species rotation and transport
 - Motivation: to understand heavy ion (W, Ar, or Ne) transport
 - Findings
 - Toroidal torque balance different for different ion species
 - Radial fluxes of different ion species can be opposite
 - Probably due to different collisionality regime distributions



- This research effort is in preliminary stage and need some guidance
 - A dedicated talk slides in Appendix B

CONCLUSIONS

- Theoretical improvement of ERT for the edge rotation study will continue
- Computational code development will continue
- Publications to follow to report the progress and findings to the plasma physics community
- Any questions or collaboration discussions can be emailed to my permanent email below
 - yuri157@gmail.com

■ Thank you for your attention!

■ Questions and Answers

Thank you for your attention!

my permanent contact: yuri157@gmail.com

APPENDIX A: GENERALIZED VISCOUS EFFECTS FOR NON-AXISYMMETRIC TOKAMAK PLASMAS

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January 2017

KSTAR conference, Muju resort, South Korea



Contents

- 1. NEOCLASSICAL VISCOUS TORQUES**
- 2. WHERE AM I TODAY?**
- 3. WHERE AM I GOING NEXT?**
- 4. WHAT CAN THIS WORK DO FOR US?**
- 5. QUESTIONS/DISCUSSIONS**

NEOCLASSICAL VISCOUS TORQUES

- Viscous forces in tokamak coordinates
 - Toroidal and poloidal momentum balance equations

$$n_j m_j \left(\vec{V}_j \cdot \nabla \right) \vec{V}_{j\phi} + \left(\nabla \cdot \vec{\pi}_j \right)_\phi = n_j e_j E_\phi^A + n_j e_j V_{rj} B_\theta + F_{\phi j} + M_{\phi j} - n_j m_j V_{atomj} V_{\phi j}$$

$$n_j m_j \left(\vec{V}_j \cdot \nabla \right) \vec{V}_{j\theta} + \frac{1}{h_\theta} \frac{\partial p_j}{\partial \theta} + \left(\nabla \cdot \vec{\pi}_j \right)_\theta = F_{\theta j} + n_j e_j E_\theta - n_j e_j V_{rj} B_\phi + S_{\theta j}^1 - m_j V_{\theta j} S_j^0$$

- Viscosity evolution equation

$$-\Omega_a \left(\vec{\pi}_a \times \hat{b} - \hat{b} \times \vec{\pi}_a \right) = \underbrace{2 p_a \nabla \vec{V}_a}_{\text{Braginskii}} + \underbrace{\frac{4}{5} \nabla q_a}_{\text{Mikhailovskii-Tsypin}} + \cancel{\partial_i \vec{\pi}_a + \left(\vec{V}_a \cdot \nabla \vec{\pi}_a + \nabla \cdot \vec{V}_a \vec{\pi}_a \right) + 2 \vec{\pi}_a \cdot \nabla \vec{V}_a + \nabla \cdot \vec{\sigma}_a - \vec{C}_a^{20}}$$

- Braginskii's viscosity representation

$$\vec{\pi}_{\alpha\beta} = \underbrace{\left(-\eta_0 W_{\alpha\beta}^0 \right)}_{\pi_{\alpha\beta}^0} + \underbrace{\left(\eta_3 W_{\alpha\beta}^3 + \eta_4 W_{\alpha\beta}^4 \right)}_{\pi_{\alpha\beta}^{34}} - \underbrace{\left(\eta_1 W_{\alpha\beta}^1 + \eta_2 W_{\alpha\beta}^2 \right)}_{\pi_{\alpha\beta}^{12}} = \underbrace{\pi_{\alpha\beta}^0}_{\parallel \text{ viscosity}} + \underbrace{\pi_{\alpha\beta}^{34}}_{\text{gyroviscosity}} + \cancel{\underbrace{\pi_{\alpha\beta}^{12}}_{\perp \text{ viscosity}}}$$

where

$$\underbrace{\eta_0 = 0.96 n T \tau}_{\parallel} \gg \underbrace{\eta_3 = \frac{1}{2} \frac{n T}{\Omega}}_{g^v}, \quad \eta_4 = 2 \eta_3 \gg \underbrace{\eta_1 = \frac{3}{10} \frac{n T}{\Omega^2 \tau}}_{\perp}, \quad \eta_2 = 4 \eta_1 \Rightarrow \left(\underbrace{\eta_0}_{\parallel} \gg \underbrace{\eta_{3,4}}_{g^v} \gg \underbrace{\eta_{1,2}}_{\perp} \right)$$

What do I mean by “Generalized Viscous Effects”?

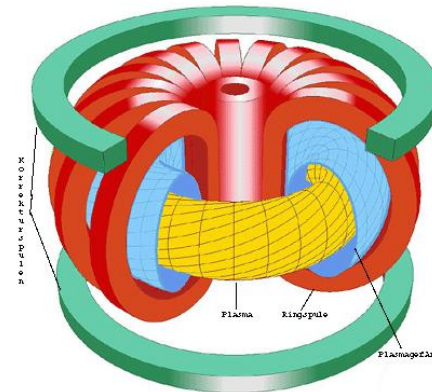
Neoclassical Toroidal Viscous (NTV) torque

$$n_j m_j R \left(\vec{V}_j \cdot \nabla \right) \vec{V}_{j\phi} + \boxed{R \left(\nabla \cdot \vec{\pi}_j \right)_\phi} = R n_j e_j E_\phi^A + R n_j e_j V_{rj} B_\theta + R F_{\phi j} + R M_{\phi j} - R n_j m_j v_{atomj} V_{\phi j}$$

$$R^2 \nabla \phi \cdot \left(\nabla \cdot \vec{\pi} \right) = \frac{1}{R h_p} \frac{\partial \left(R^2 h_p \pi_{\psi\phi} \right)}{\partial l_\psi} + B_p \frac{\partial}{\partial l_p} \left(\frac{R \pi_{p\phi}}{B_p} \right) + \frac{\partial \pi_{\phi\phi}}{\partial l_\phi}$$

↓ take Flux Surface Average

$$\left\langle R^2 \nabla \phi \cdot \left(\nabla \cdot \vec{\pi} \right) \right\rangle = \left\langle \frac{1}{R h_p} \frac{\partial \left(R^2 h_p \pi_{\psi\phi} \right)}{\partial l_\psi} \right\rangle + \left\langle B_p \frac{\partial}{\partial l_p} \left(\frac{R \pi_{p\phi}}{B_p} \right) \right\rangle + \left\langle \frac{\partial \pi_{\phi\phi}}{\partial l_\phi} \right\rangle$$



Neoclassical Poloidal Viscous (NPV) torque

$$r n_j m_j \left(\vec{V}_j \cdot \nabla \right) \vec{V}_{j\theta} + \frac{r}{h_\theta} \frac{\partial p_j}{\partial \theta} + \boxed{r \left(\nabla \cdot \vec{\pi}_j \right)_\theta} = r F_{\theta j} + r n_j e_j E_\theta - r n_j e_j V_{rj} B_\phi + r S_{\theta j}^1 - r m_j V_{\theta j} S_j^0$$

$$\left(\nabla \cdot \vec{\pi} \right)_p = \sum_{j=\psi, p, \phi} \left[\frac{1}{h_\psi h_p h_\phi} \frac{\partial}{\partial j} \left(\frac{h_\psi h_p h_\phi}{h_j} \pi_{jp} \right) + \sum_{k=\psi, p, \phi} \Gamma_{pk}^j \pi_{jk} \right]$$

Generalized viscous effects

- Work out NTV and NPV torques in non-axisymmetric plasmas

Flux Surface Averaged NTV torque in non-axisymmetric plasmas

$$n_j m_j \left\langle \left[R \left(\vec{V}_j \cdot \nabla \right) \vec{V}_j \right]_{\phi} \right\rangle + \boxed{\left\langle \left[R \left(\nabla \cdot \vec{\pi}_j \right) \right]_{\phi} \right\rangle} = \langle R n_j e_j E_{\phi}^A \rangle + \langle R n_j e_j V_{rj} B_{\theta} \rangle + \langle R F_{\phi j} \rangle + \langle R M_{\phi j} \rangle - \langle R n_j m_j v_{atomj} V_{\phi j} \rangle$$

NTV Torque

$$\left\langle R^2 \nabla \phi \cdot (\nabla \cdot \vec{\pi}) \right\rangle = \left\langle \frac{1}{R h_p} \frac{\partial (R^2 h_p \pi_{\psi\phi})}{\partial l_{\psi}} \right\rangle$$

$$\begin{matrix} \eta_0 \gg \eta_{3,4} \gg \eta_{1,2} \\ \parallel \quad \quad \quad \perp \\ g v \quad \quad \quad \end{matrix}$$

$$\left\langle R^2 \nabla \phi \cdot (\nabla \cdot \vec{\pi}) \right\rangle = \underbrace{\left\langle \frac{1}{R h_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^0)}{\partial l_{\psi}} \right\rangle}_{\text{non-axisymmetric}} + \underbrace{\left\langle \frac{1}{R h_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{34})}{\partial l_{\psi}} \right\rangle}_{\text{axisymmetric}} + \cancel{\left\langle \frac{1}{R h_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{12})}{\partial l_{\psi}} \right\rangle}$$

$$\pi_{\psi\phi} \approx \pi_{\psi\phi}^0 + \pi_{\psi\phi}^{34} = -\eta_0 W_{\psi\phi}^0 + \eta_4 \left(\frac{1}{2} W_{\psi\phi}^3 + W_{\psi\phi}^4 \right) - \cancel{\eta_2 \left(\frac{1}{4} W_{\psi\phi}^1 + W_{\psi\phi}^2 \right)}$$

$$W_{\psi\phi}^0 = \frac{3}{2} \underbrace{f_{\psi}}_{\text{blue dashed}} f_p H^0$$

$$f_{\psi} \equiv \frac{B_{\psi}}{B}, \quad f_p \equiv \frac{B_p}{B}, \quad f_{\phi} \equiv \frac{B_{\phi}}{B}$$

$$H^0 = \sum_{\mu, \nu} \left(f_{\mu} f_{\nu} - \frac{1}{3} \delta_{\mu\nu} \right) W_{\mu\nu}$$

$$W_{\mu\nu} = \underbrace{\left(\overline{\nabla \vec{V}} \right)_{\mu\nu}}_{\text{Braginskii}} + \underbrace{\left(\frac{2}{5p} \nabla \vec{q} \right)_{\mu\nu}}_{\text{Mikhailovskii-Tsypin}}$$

WHERE AM I TODAY?

- **Axisymmetric** NTV and NPV torques calculated in GTROTA
 - NTV represented with gyroviscosity / Braginskii's closure

$$\left\langle R^2 \nabla \phi \cdot (\nabla \cdot \tilde{\pi}) \right\rangle = \underbrace{\left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^0)}{\partial l_\psi} \right\rangle}_{\text{non-axisymmetric}} + \underbrace{\left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{34})}{\partial l_\psi} \right\rangle}_{\text{axisymmetric}} + \left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{12})}{\partial l_\psi} \right\rangle$$

$$W_{\mu\nu} = \underbrace{\left(\overline{\nabla \nabla} \right)_{\mu\nu}}_{\text{valid for high collisionality}}^{\text{Braginskii}} + \left(\frac{2}{5p} \overline{\nabla q} \right)_{\mu\nu}$$

❖ Two neoclassical codes that handle gyroviscosity: GTROTA / NEO

- **GTROTA based on collisionality-extended Braginskii's viscosity**
 - Parallel viscosity coefficient extended to low collisionality (trapped particle effect) by Shaing

$$\underbrace{\eta_0 = 0.96 n T \tau}_{\text{collisionality dependent}}, \quad \underbrace{\eta_3 = \frac{1}{2} \frac{n T}{\Omega}}_{\text{independent of collisionality}}, \quad \underbrace{\eta_4 = 2 \eta_3}_{\text{independent of collisionality}}, \quad \underbrace{\eta_1 = \frac{3}{10} \frac{n T}{\Omega^2 \tau}}_{\text{collisionality dependent (but small contribution)}}, \quad \underbrace{\eta_2 = 4 \eta_1}_{\text{collisionality dependent (but small contribution)}}, \quad \Rightarrow \quad \left(\eta_0 \gg \eta_{3,4} \gg \eta_{1,2} \right)$$

■ Shaing-banana plateau-PS:

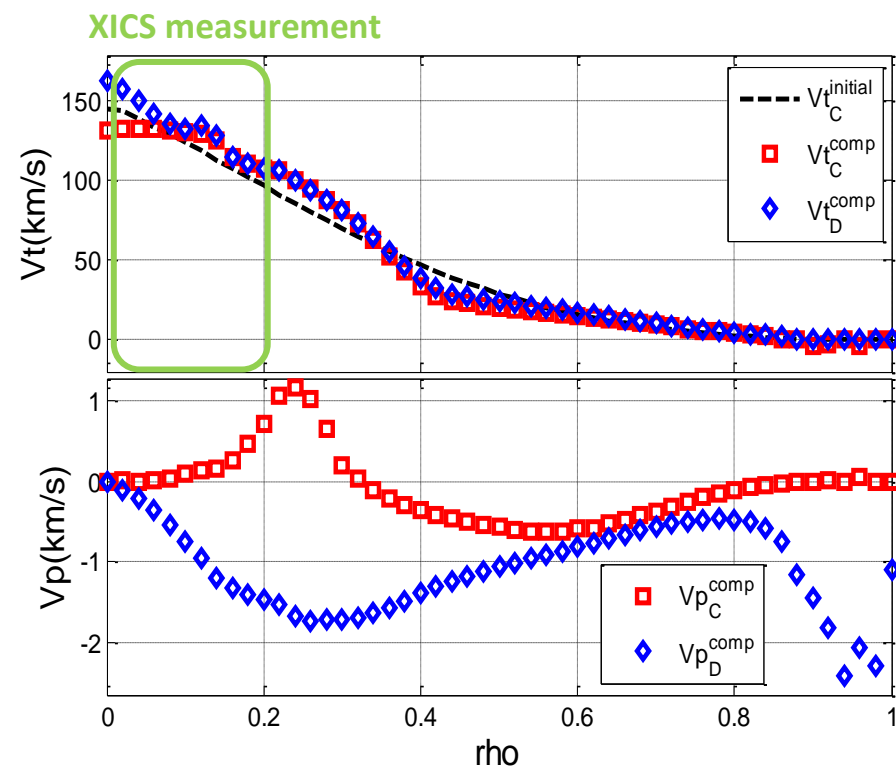
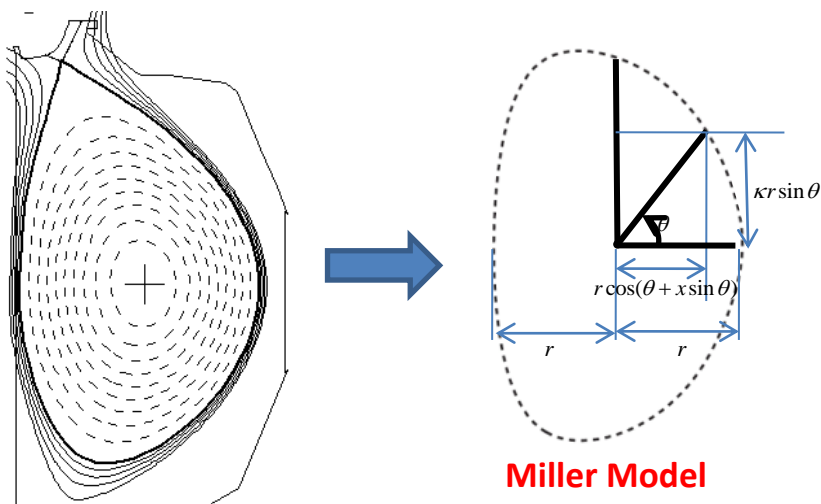
$$\eta_{0j} = \frac{n_j m_j V_{thj} q R_0 \varepsilon^{-3/2} v_{jj}^*}{(1 + \varepsilon^{-3/2} v_{jj}^*)(1 + v_{jj}^*)} \equiv n_j m_j V_{thj} q R f_j$$

-> **valid for arbitrary collisionality**

- GTROTA considers first-order poloidal asymmetries in density, velocity, and electrostatic potential
 - Rotation calculations with KSTAR #5505-2500ms (H-mode with NBI)
 - [Bae et al., PoP 21, 012504(2014)]
 - No non-axisymmetric magnetic perturbation in this shot

$$\left\langle R^2 \nabla \phi \cdot (\nabla \cdot \vec{\pi}) \right\rangle = \underbrace{\left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^0)}{\partial l_\psi} \right\rangle}_{\text{non-axisymmetric}} + \underbrace{\left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{34})}{\partial l_\psi} \right\rangle}_{\text{axisymmetric}} + \left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{12})}{\partial l_\psi} \right\rangle$$

$$\pi_{\psi\phi} \approx \pi_{\psi\phi}^0 + \pi_{\psi\phi}^{34} = -\eta_0 W_{\psi\phi}^0 + \eta_4 \left(\frac{1}{2} W_{\psi\phi}^3 + W_{\psi\phi}^4 \right) - \eta_2 \left(\frac{1}{4} W_{\psi\phi}^1 + W_{\psi\phi}^2 \right)$$



WHERE AM I GOING NEXT?

- Generalize to **non-axisymmetric** NTV and NPV torques
 - Theoretical study published [Stacey and Bae, PoP 2015]
 - Theoretical model development for GTROTA in progress
 - Extend to Mikhailovskii-Tsypin's closure
 - need Heat Flux Density Evolution Equation

$$\left\langle R^2 \nabla \phi \cdot (\nabla \cdot \vec{\pi}) \right\rangle = \underbrace{\left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^0)}{\partial l_\psi} \right\rangle}_{\text{non-axisymmetric}} + \underbrace{\left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{34})}{\partial l_\psi} \right\rangle}_{\text{axisymmetric}} + \left\langle \frac{1}{Rh_p} \frac{\partial (R^2 h_p \pi_{\psi\phi}^{12})}{\partial l_\psi} \right\rangle$$

$$\pi_{\psi\phi} \approx \pi_{\psi\phi}^0 + \pi_{\psi\phi}^{34} = -\eta_0 W_{\psi\phi}^0 + \eta_4 \left(\frac{1}{2} W_{\psi\phi}^3 + W_{\psi\phi}^4 \right) - \eta_2 \left(\frac{1}{4} W_{\psi\phi}^1 + W_{\psi\phi}^2 \right)$$

$$W_{\psi\phi}^0 = \frac{3}{2} f_\psi f_p H^0$$

$$W_{\mu\nu} = \overbrace{\left(\vec{\nabla} \vec{V} \right)_{\mu\nu}}^{\text{Braginskii}} + \left(\frac{2}{5p} \vec{\nabla} \vec{q} \right)_{\mu\nu}$$

$$-\Omega_a (\hat{b} \times \vec{q}_a) = \frac{5}{2m_a} p_a \nabla T_a + \frac{1}{2} \vec{C}_a^{11} + O(\text{higher})$$

$$H^0 \equiv \left[\begin{aligned} & \left(f_\psi f_\psi - \frac{1}{3} \right) \left\{ \frac{4}{3} \frac{\partial V_\psi}{\partial l_\psi} - \frac{2}{3} \left(\frac{\partial V_p}{\partial l_\psi} + \frac{\partial V_\phi}{\partial l_\psi} \right) + 2 \left(\frac{1}{h_\psi} \frac{\partial h_\psi}{\partial l_p} V_p + \frac{1}{h_\psi} \frac{\partial h_\psi}{\partial l_\phi} V_\phi \right) \right\} + \\ & \left(f_p f_p - \frac{1}{3} \right) \left\{ \frac{4}{3} \frac{\partial V_p}{\partial l_p} - \frac{2}{3} \left(\frac{\partial V_\psi}{\partial l_p} + \frac{\partial V_\phi}{\partial l_p} \right) + 2 \left(\frac{1}{h_p} \frac{\partial h_p}{\partial l_\psi} V_\psi + \frac{1}{h_p} \frac{\partial h_p}{\partial l_\phi} V_\phi \right) \right\} + \\ & \left(f_\phi f_\phi - \frac{1}{3} \right) \left\{ \frac{4}{3} \frac{\partial V_\phi}{\partial l_\phi} - \frac{2}{3} \left(\frac{\partial V_p}{\partial l_\phi} + \frac{\partial V_\psi}{\partial l_\phi} \right) + 2 \left(\frac{1}{h_\phi} \frac{\partial h_\phi}{\partial l_\psi} V_\psi + \frac{1}{h_\phi} \frac{\partial h_\phi}{\partial l_p} V_p \right) \right\} + \\ & 2f_\psi f_p \left\{ \frac{\partial V_\psi}{\partial l_p} + \frac{\partial V_p}{\partial l_\psi} - \left(\frac{1}{h_p} \frac{\partial h_p}{\partial l_\psi} V_p + \frac{1}{h_\psi} \frac{\partial h_\psi}{\partial l_p} V_\psi \right) \right\} + \\ & 2f_p f_\phi \left\{ \frac{\partial V_\phi}{\partial l_p} + \frac{\partial V_p}{\partial l_\phi} - \left(\frac{1}{h_p} \frac{\partial h_p}{\partial l_\phi} V_p + \frac{1}{h_\phi} \frac{\partial h_\phi}{\partial l_p} V_\phi \right) \right\} + \\ & 2f_\psi f_\phi \left\{ \frac{\partial V_\phi}{\partial l_\psi} + \frac{\partial V_\psi}{\partial l_\phi} - \left(\frac{1}{h_\psi} \frac{\partial h_\psi}{\partial l_\phi} V_\psi + \frac{1}{h_\phi} \frac{\partial h_\phi}{\partial l_\psi} V_\phi \right) \right\} \end{aligned} \right]$$

WHAT CAN THIS WORK DO FOR US?

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■ Rotation and transport predictions of deuterium for non-axisymmetric plasmas

■ Generalized viscous contribution to rotation and transport

$$\left\langle n_j m_j \left[R(\vec{V}_j \cdot \nabla) \vec{V}_j \right]_{\phi} \right\rangle + \left\langle \left[R(\nabla \cdot \vec{\pi}_j) \right]_{\phi} \right\rangle = \left\langle R n_j e_j E_{\phi}^A \right\rangle + \left\langle R n_j e_j V_{rj} B_{\theta} \right\rangle + \left\langle R F_{\phi j} \right\rangle + \left\langle R M_{\phi j} \right\rangle - \left\langle R n_j m_j v_{atomj} V_{\phi j} \right\rangle$$

$$\left\langle m_j m_j (\vec{V}_j \cdot \nabla) \vec{V}_{j\theta} \right\rangle + \left\langle \frac{r}{h_{\theta}} \frac{\partial p_j}{\partial \theta} \right\rangle - \left\langle r(\nabla \cdot \vec{\pi}_j)_{\theta} \right\rangle = \left\langle r F_{\theta j} \right\rangle + \left\langle r n_j e_j E_{\theta} \right\rangle - \left\langle r n_j e_j V_{rj} B_{\phi} \right\rangle + \left\langle r S_{\theta j}^1 \right\rangle - \left\langle r m_j V_{\theta j} S_j^0 \right\rangle$$



$$\vec{\Gamma}_j = n_j \vec{V}_j$$

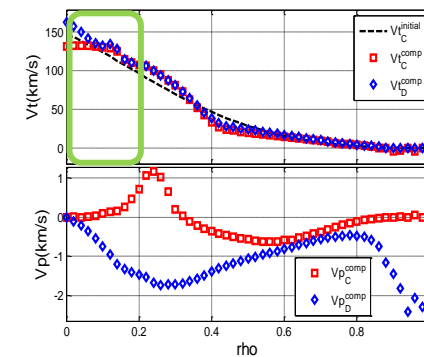
$$\vec{M}_j$$

$$\vec{q}_j$$



KSTAR #5505

XICS measurement



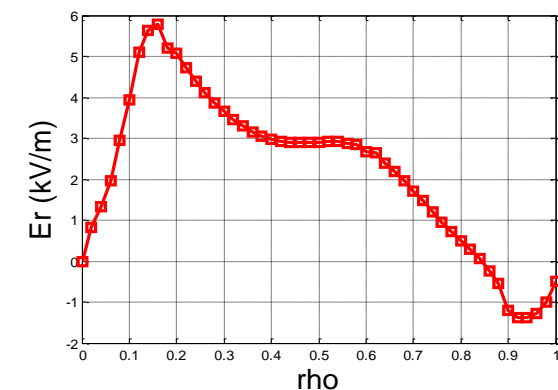
■ Increase accuracy in Er calculation

$$\bar{E}_r = -V_{thj} \bar{B}_{\theta} \left[V_{\theta j} \frac{\left\langle \frac{1}{1 + \varepsilon \cos \xi} \right\rangle}{\left\langle \frac{1}{h_r} \right\rangle} - V_{\phi j} \left(1 + \frac{\partial R_0(r)}{\partial r} \right) \frac{\left\langle \frac{1}{(1 + \varepsilon \cos \xi)} \frac{1}{h_r} \right\rangle}{\left\langle \frac{1}{h_r} \right\rangle} - \frac{1}{V_{thj}} \frac{1}{n_j e_j \bar{B}_{\theta}} \frac{\partial \bar{P}_j}{\partial r} \right]$$

$$\bar{E}_r^{cir} = \frac{1}{n_j e_j} \frac{\partial \bar{P}_j}{\partial r} - \left[V_{\theta j} \bar{B}_{\phi} - V_{\phi j} \bar{B}_{\theta} \right]$$

KSTAR #5505

Radial Electric Field





APPENDIX B:

A COUPLED STUDY OF

PLASMA ROTATION AND TRANSPORT

: COMPARISON OF TOROIDAL TORQUE CONTRIBUTIONS

IN AXISYMMETRIC TOKAMAK PLASMAS

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2. University of Science and Technology, Daejeon, South Korea

11 July 2017
Sapporo, Japan

INTRODUCTION

■ Practical questions!!!

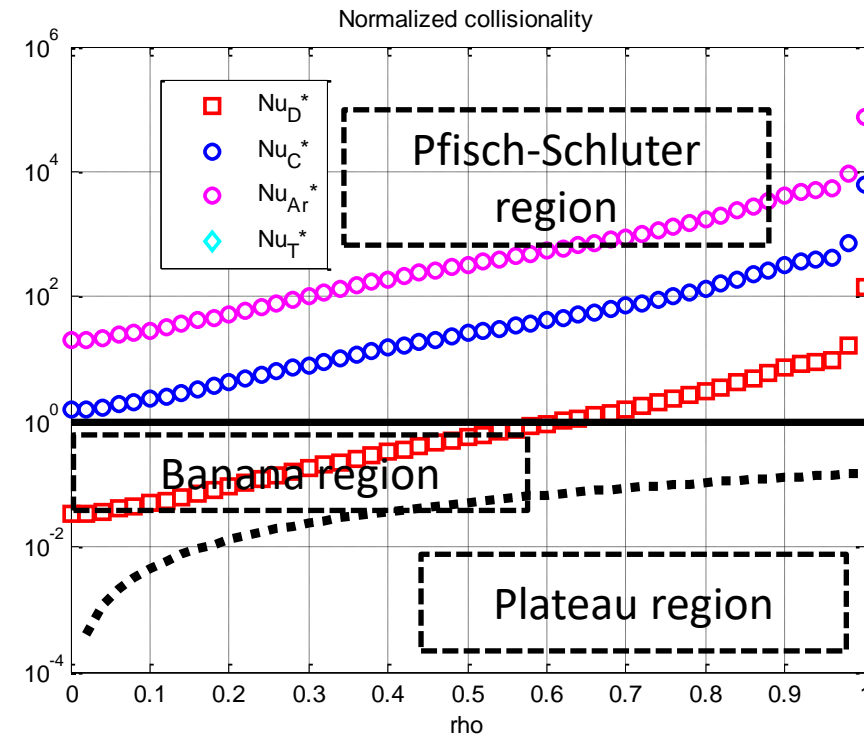
1. How are the rotation & transport differ for different ion species?

- considering their collisionalities differ significantly

2. Can we predict deuterium rotation & transport from the impurity measurements?

- Most transport studies done with impurities

3. Which physical term(s) in the angular momentum(or torque) balance equation have the largest contributions to rotation & transport?



$$\left\langle R m_a n_a \frac{\partial V_{a\phi}}{\partial t} \right\rangle + \left\langle R n_a m_a \left(\vec{V}_a \cdot \nabla \right) V_{a\phi} \right\rangle + \left\langle R \frac{1}{h_\phi} \frac{\partial p_a}{\partial \phi} \right\rangle + \left\langle R \left(\nabla \cdot \vec{\pi}_a \right)_\phi \right\rangle = \left\langle R n_a e_a \left(-\frac{1}{h_\phi} \frac{\partial \Phi}{\partial \phi} - \frac{\partial A_\phi}{\partial t} \right) \right\rangle + \left\langle R n_a e_a \left(V_{ar} B_\theta - V_{a\theta} B_r \right) \right\rangle + \left\langle R F_{a\phi}^1 \right\rangle + \left\langle R S_{a\phi}^1 \right\rangle - \left\langle R m_a V_{a\phi} S_a^0 \right\rangle$$

■ A coupled study of plasma rotation & transport to answer these questions!

INTRODUCTION

- Most common “particle (or flux)” transport study of today
 - is based on the continuity equation

$$\frac{\partial n_a}{\partial t} + \underbrace{\nabla \cdot (n_a \vec{V}_a)}_{\text{Dfn}(n_a) + \text{Vfn}(\vec{V}_a)} = S_a^0$$

- closes this equation with approximated physical effects
 - with density gradient: introduces “Diffusion coefficient (D)”
 - with convective effect: introduces “Convective coefficient (V)”
- We use this model to study individual ion transport
 - by injecting a specific impurities (Ar, Ne, W, etc.)
 - measure fluxes to determine D, V, and effective Vr
 - require dedicated injection systems and diagnostics
 - but still study transport of only one impurity
- Q: can’t we just solve the momentum balance equation(MBE) to get Vr?
 - as we do in plasma rotation study

$$n_a m_a \frac{\partial \vec{V}_a}{\partial t} + n_a m_a (\vec{V}_a \cdot \nabla) \vec{V}_a + \nabla p_a + \nabla \cdot \vec{\pi}_a = n_a e_a (\vec{E} + \vec{V}_a \times \vec{B}) + \vec{F}_a^1 + \left(\vec{S}^1 - m_a \vec{V}_a S^0 \right)$$

INTRODUCTION

- **Plasma rotation study** solves MBE to get velocities (**mostly Vt and Vp**)

$$n_a m_a \frac{\partial \vec{V}_a}{\partial t} + n_a m_a \left(\vec{V}_a \cdot \nabla \right) \vec{V}_a + \nabla P_a + \nabla \cdot \vec{\pi}_a = n_a e_a \left(\vec{E} + \vec{V}_a \times \vec{B} \right) + \vec{F}_a^1 + \left(\vec{S}^1 - m_a \vec{V}_a S^0 \right)$$

- if **Vr** can also be calculated, two studies can be coupled
- **Plasma rotation theories** are developed by two main approaches
 - **NCLASS** (based on Hirshman-Sigmar approach) [Houlberg et. al., 1998]
 - most famous with two (parallel and perpendicular) MBEs to calculate neoclassical rotations of multi-ions
 - but no further development to couple it with particle/heat transport
 - **GTROTA** (based on Stacey-Sigmar approach) [Bae et. al., Comp. Phys. Comm. 2013]
 - introduced as “**Extended Plasma Rotation Theory (EPRT)**” [Bae et. el., NF 2013]
 - takes **MBE in three coordinates (radial, poloidal, toroidal)**
 - direct comparison with **Vt and Vp measurements** possible
 - **Vr** specifically appears in the toroidal torque balance (next slide)
- **Q: can I extend EPRT/GTROTA to find Vr/radial fluxes for most generalized tokamak plasmas(both axisymmetric and non-axisymmetric)?**

CURRENT STATUS OF EPRT/GTROTA

- Retains **all the terms** in the MBEs and in **three** coordinates
- Velocity calculation models** for **Vt** and **Vp**
 - Toroidal direction: **toroidal torque balance** (Flux Surface Averaged)
 - Vr** appears in this equation => **Effective Vr** (later slide)

$$\left\langle Rm_a n_a \frac{\partial V_{a\phi}}{\partial t} \right\rangle + \left\langle Rn_a m_a \left(\vec{V}_a \cdot \nabla \right) V_{a\phi} \right\rangle + \left\langle R \frac{1}{h_\phi} \frac{\partial p_a}{\partial \phi} \right\rangle + \left\langle R \left(\nabla \cdot \vec{\pi}_a \right)_\phi \right\rangle =$$

$$\left\langle Rn_a e_a \left(-\frac{1}{h_\phi} \frac{\partial \Phi}{\partial \phi} - \frac{\partial A_\phi}{\partial t} \right) \right\rangle + \left\langle Rn_a e_a \left(V_{ar} B_\theta - V_{a\theta} B_r \right) \right\rangle + \left\langle RF_{a\phi}^1 \right\rangle + \left\langle RS_{a\phi}^1 \right\rangle - \left\langle Rm_a V_{a\phi} S_a^0 \right\rangle$$

- Poloidal direction: **poloidal torque balance** (Flux Surface Averaged)

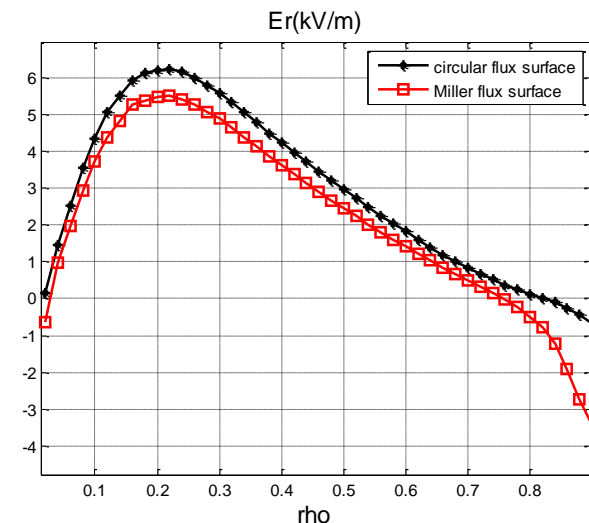
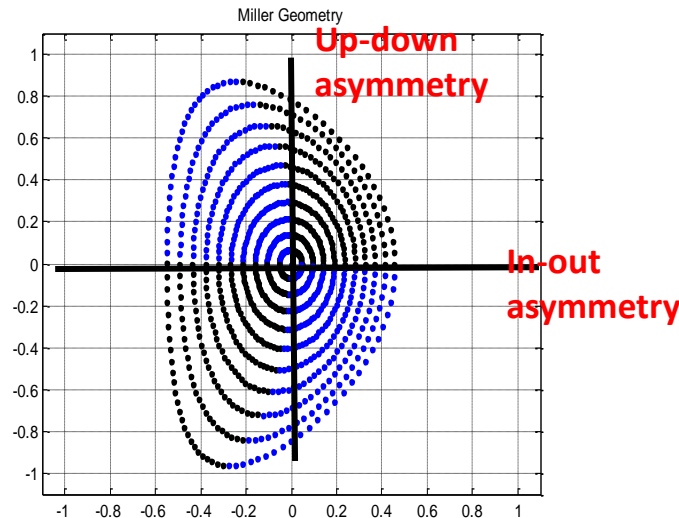
$$\left\langle rm_a n_a \frac{\partial V_{a\theta}}{\partial t} \right\rangle + \left\langle rn_a m_a \left(\vec{V}_a \cdot \nabla \right) V_{a\theta} \right\rangle + \left\langle r \frac{1}{h_\theta} \frac{\partial p_a}{\partial \theta} \right\rangle + \left\langle r \left[\nabla \cdot \vec{\pi}_a \right]_\theta \right\rangle =$$

$$\left\langle r \left(-\frac{1}{h_p} \frac{\partial \Phi}{\partial \theta} - \frac{\partial A_\theta}{\partial t} \right) \right\rangle + \left\langle rn_a e_a \left(V_{a\phi} B_r - V_{ar} B_\phi \right) \right\rangle + \left\langle rF_{a\theta}^1 \right\rangle + \left\langle r \left(S_{a\theta}^1 - m_a V_{a\theta} S_a^0 \right) \right\rangle$$

- Radial direction: **radial MBE**
 - Provides coupling relations with **the continuity equation**
 - Used to calculate 1st order poloidal variations (aka **poloidal asymmetries**)
 - in **density, velocity, and electrostatic potential** [Bae et. Al., PoP 2014]

CURRENT STATUS OF EPRT/GTROTA

- Other EPRT/GTROTA features
 - Uses **Miller flux surface** geometry
 - for higher accuracy in **velocities**(V_t , V_p , V_r), **momentum/particle transport**, and **E_r** calculations



- All** plasma parameters (V_t , V_p , V_r , E_r , etc.) **self-consistently** iterated
 - Maximizes the advantages of plasma fluid equations
- Calculates **E_r** , **V_t** , **V_p** , **V_r** , and **poloidal asymmetries** of up to **four ion species & electron**
 - Simulation of multi-ion plasmas possible (eg., D-T plasma with Tungsten)

CURRENT STATUS OF EPRT/GTROTA

- Current GTROTA handles **axisymmetric** plasmas only
 - Non-axisymmetric theory available [Stacey and Bae, PoP 2015] but not in GTROTA
- Today, calculated **Effective Vr** in **axisymmetric** plasmas
 - defined to represent **all anomalous terms** (2nd order and higher / turbulence inclusive)
 - Assuming most **anomalous** transports are in **radial** direction

Accelerational

Convective

Pressure

Viscous (NTV)

$$\left\langle Rm_a n_a \frac{\partial V_{a\phi}}{\partial t} \right\rangle + \left\langle Rn_a m_a (\vec{V}_a \cdot \nabla) V_{a\phi} \right\rangle + \left\langle R \frac{1}{h_\phi} \frac{\partial p_a}{\partial \phi} \right\rangle + \left\langle R (\nabla \cdot \vec{\pi}_a)_\phi \right\rangle =$$

$$\left\langle Rn_a e_a \left(-\frac{1}{h_\phi} \frac{\partial \Phi}{\partial \phi} - \frac{\partial A_\phi}{\partial t} \right) \right\rangle + \left\langle Rn_a e_a (V_{ar} B_\theta - V_{a\theta} B_r) \right\rangle + \left\langle RF_{a\phi}^1 \right\rangle + \left\langle RS_{a\phi}^1 \right\rangle - \left\langle Rm_a V_{a\phi} S_a^0 \right\rangle$$

Electric

V cross B

Frictional External Atomic

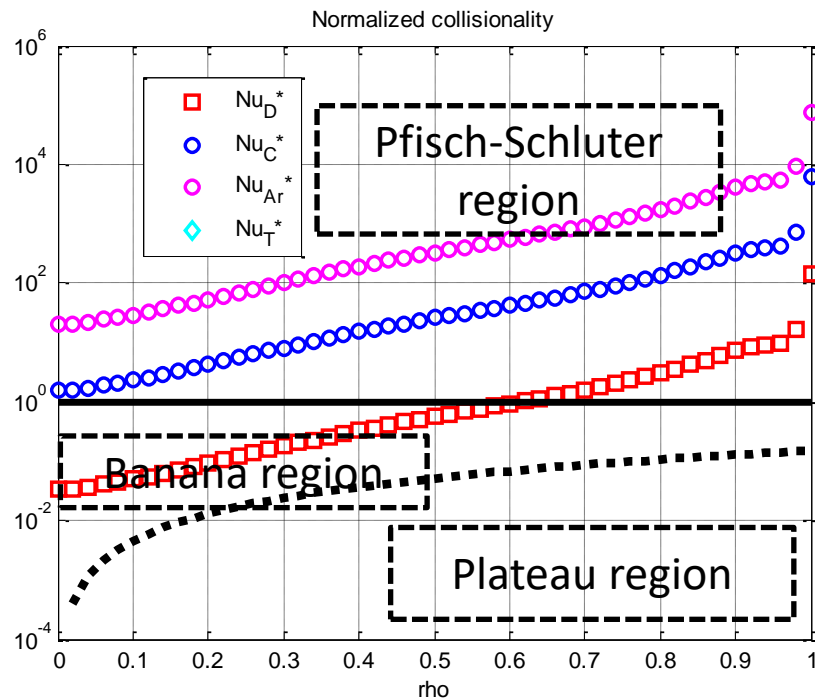


**“Anomalous
(Effective Vr cross B)”**

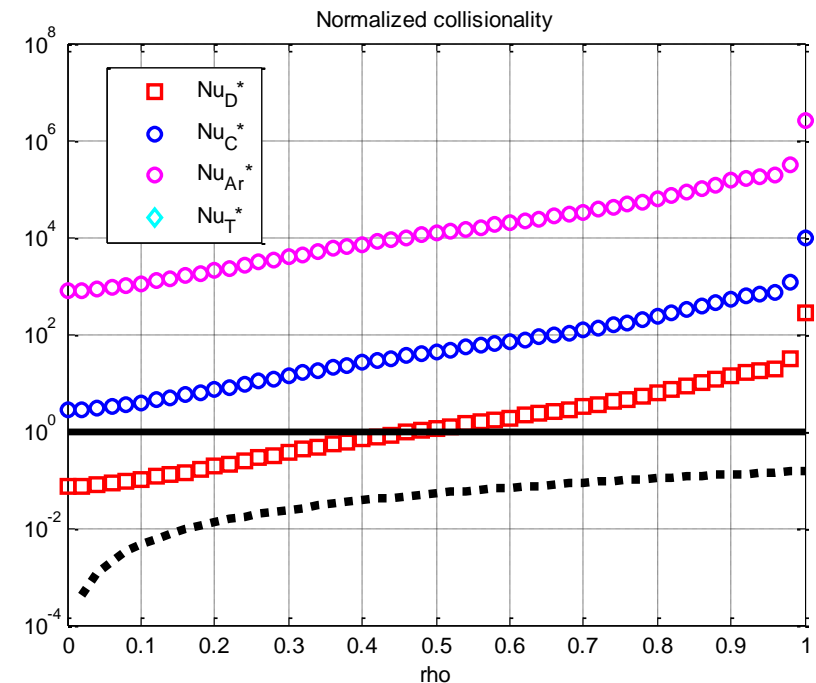
$$\underbrace{\left\langle Rn_a e_a V_{ar}^{eff} B_\theta \right\rangle}_{Rn_a e_a B_\theta (V_r^{eff})} = \text{grey terms} + \underbrace{\text{blue terms}}_{\text{Reynolds stress}}$$

DISCHARGES ANALYZED

- Four KSTAR discharges analyzed
 - Two H-modes with Ar rotation measured (#5505-2500ms / #5953-2500ms)
 - Two L-modes with Carbon rotation measured (#13728-4500ms / 13728-4950ms)
- One simulation with Tungsten
 - Based on #5505-2500ms



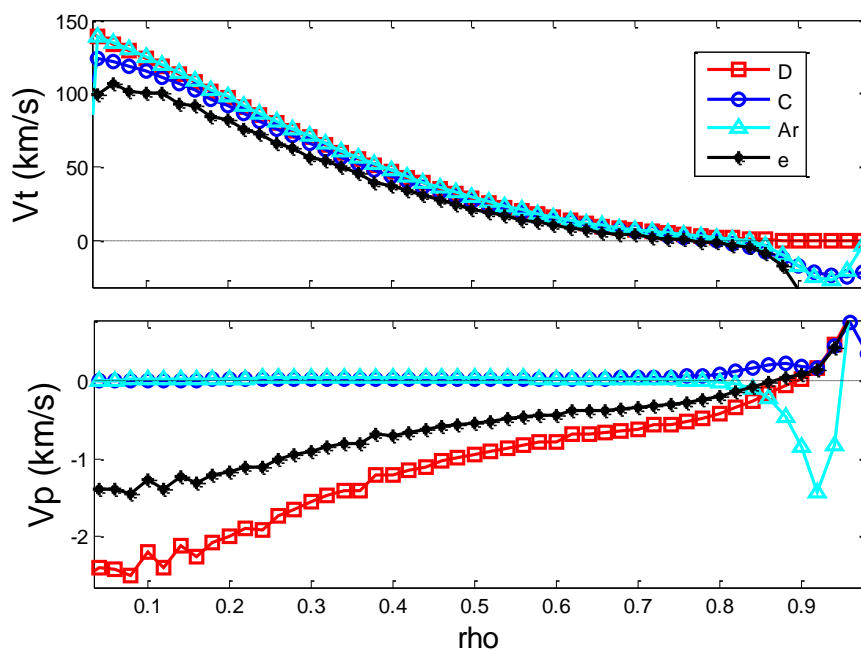
■ KSTAR discharge #5505 (2500ms)



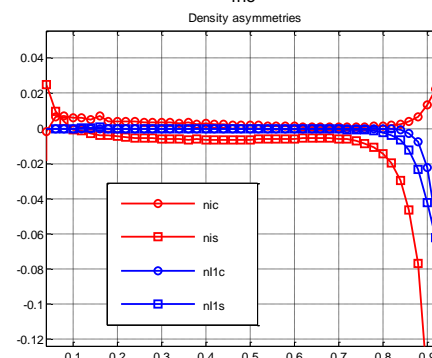
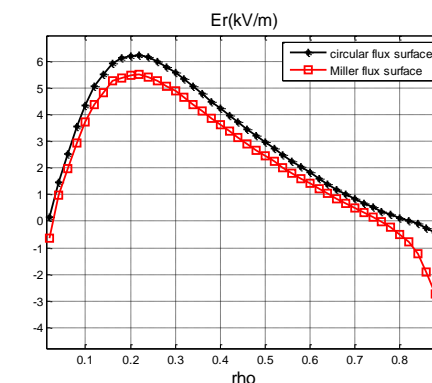
■ KSTAR discharge #5505 (2500ms) with Ar replace with W

ANALYSIS RESULTS

- **Rotation, poloidal asymmetries, & E_r (KSTAR #5505-2500ms)**
 - V_p and V_t of all ions/electrons calculated
 - V_t very close to each other and stays **within 10% of the measurement** [Bae et. et., NF 2013 / Bae et. al., PoP 2014]
 - E_r self-consistently(iteratively) calculated
 - **Poloidal asymmetries (of density, velocity, and electrostatic potential) calculated** (density asymmetry only shown in this slide)

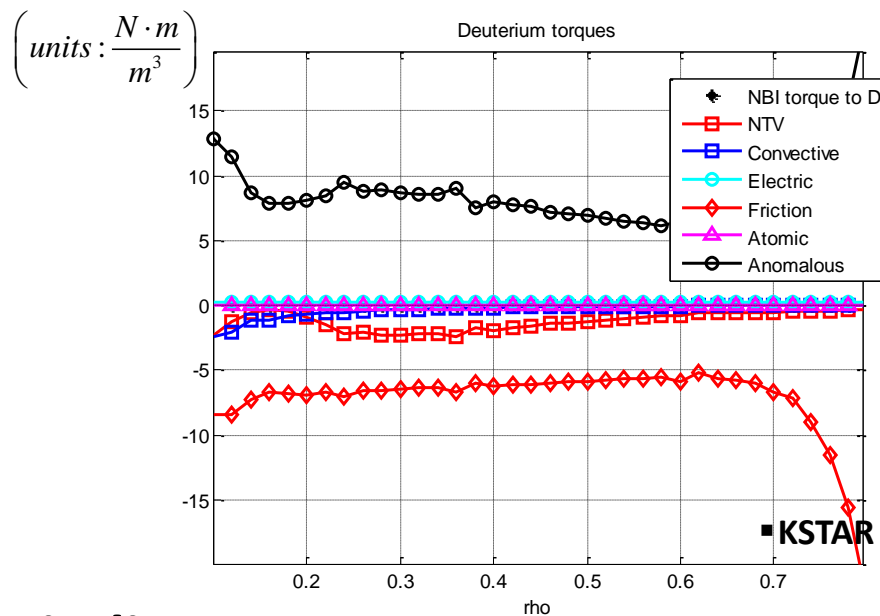


■ KSTAR discharge #5505 (2500ms)

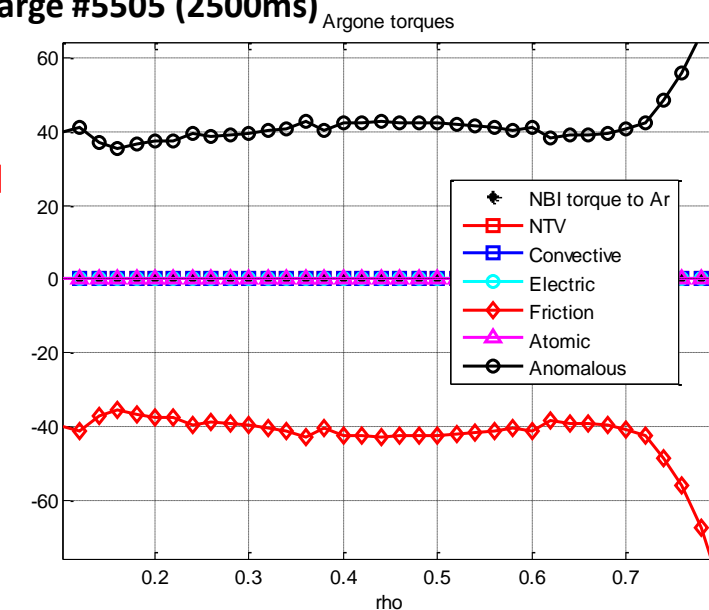
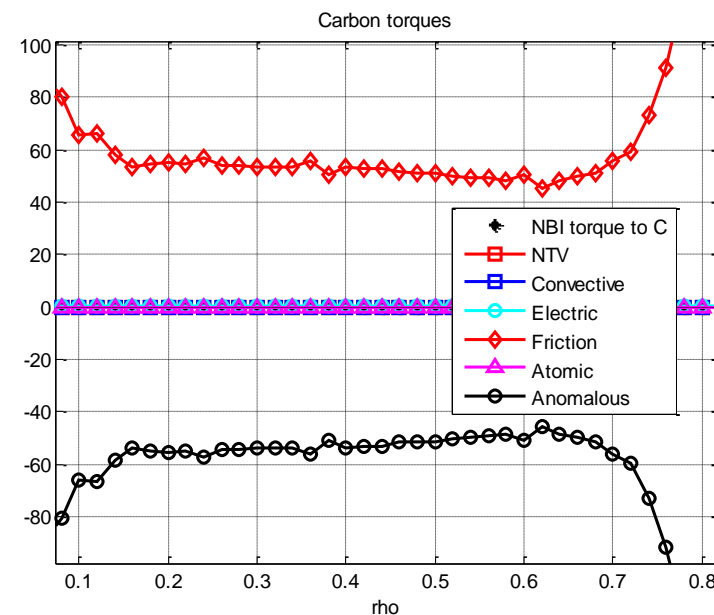


TOROIDAL TORQUE COMPARISONS

Toroidal torque densities



■ KSTAR discharge #5505 (2500ms)

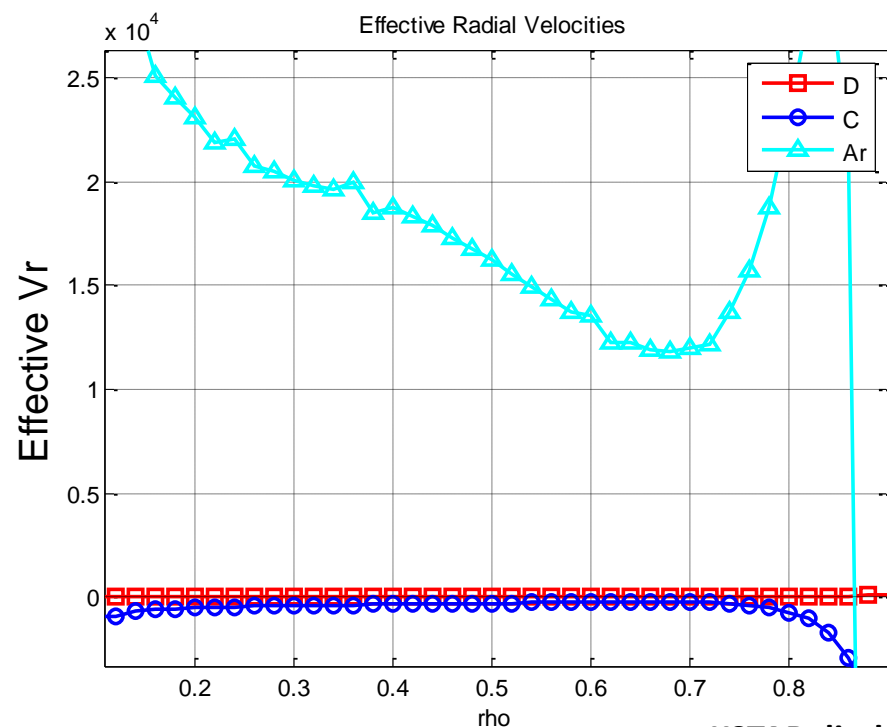


Findings

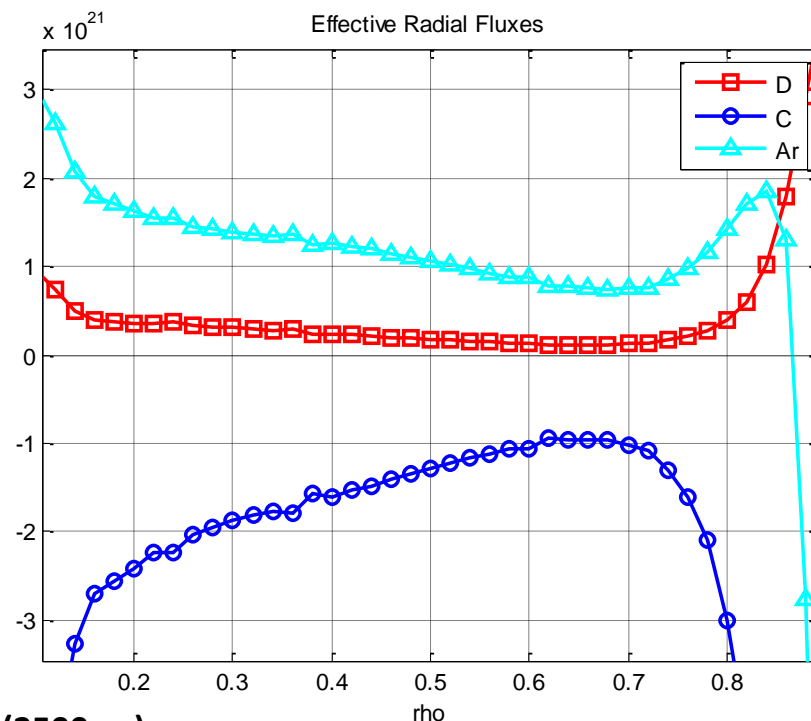
- torque balance mostly maintained by **friction and anomalous** ("Effective V_r cross B ")
 - especially for the impurities
- Different balancing relations for different ion species
 - Q: Is "collisionality" most dominant effect in the rotation and transport of different ion species?

TOROIDAL TORQUE COMPARISONS

Effective Vr / radial fluxes



■ KSTAR discharge #5505 (2500ms)



Findings

- Effective Vr & radial fluxes **larger for impurities**
- Impurity fluxes in **opposite directions**
 - Q: Do all the impurities always accumulate in the core?
 - Q: What physical mechanism determines these directions? Collisionality?

FUTURE OPPORTUNITIES

- **Theoretical model development (in progress)**
 - develop a separate **Vr subsystem**
 - develop numerical models for **non-axisymmetric** plasmas [Stacey and Bae, PoP 2015]
 - include **1st order toroidal variations (in Bt and others)**

- **Numerical code development**

- Code in **non-axisymmetric** theories
 - compare all the **poloidal** torque density terms

$$\left\langle rm_a n_a \frac{\partial V_{ap}}{\partial t} \right\rangle + \left\langle rn_a m_a (\vec{V}_a \cdot \nabla) V_{ap} \right\rangle + \left\langle r \frac{1}{h_p} \frac{\partial p_a}{\partial p} \right\rangle + \left\langle r \left[\nabla \cdot \vec{\pi}_a \right]_p \right\rangle =$$

$$\left\langle r \left(-\frac{1}{h_p} \frac{\partial \Phi}{\partial p} - \frac{\partial A_p}{\partial t} \right) \right\rangle + \left\langle rn_a e_a (V_{a\phi} B_\psi - V_{a\psi} B_\phi) \right\rangle + \left\langle r F_{ap}^1 \right\rangle + \left\langle r (S_{ap}^1 - m_a V_{ap} S_a^0) \right\rangle$$

- **Simulations/analysis of modern tokamaks**

- Analyze KSTAR discharges with **ITBs**
 - Various modes/devices: **DIID-D, KSTAR, EAST, etc.**
 - **W** transport simulations
 - **ITER**-relevant simulations (ITER shapes / D-T fusion / etc.)
 - **Collisionality** effect simulations

- **open to comments and ideas!**

Thank you for your attention!

Additional slides with more details...

1. Simulation with Tungsten
2. L-mode discharge analysis results
3. Extended Plasma Rotation theory

my permanent contact: yuri157@gmail.com

APPENDIX C:

EXPERIMENTAL EXPERIENCES WITH KSTAR

KSTAR CAMPAIGN EXPERIENCES

CPO/APO experiences [2013-1017]

- Served as CPO/APO for five consecutive KSTAR annual campaigns
- Develop CPO/APO training manual and checklist

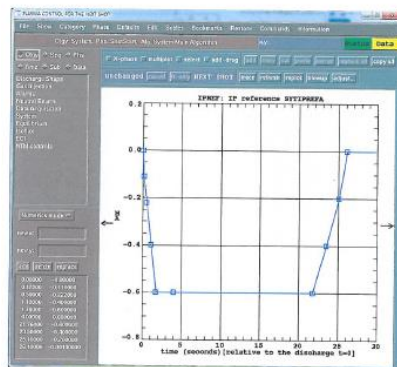
PCS/WAVE 시스템 구성개념

Ctgy(직원)-Seq(상황별 시간표)-Phase(수행업무)-Sub(업무별 세부절차)- 각 Sub 의 항목들(세부절차별 행동요령) 관계

by 배천호 (2014-11-20)

Ctgy 종류: 10 가지

* 이것은 내 회사에 같이 일을 하는 10 사람의 직원이 있는 것과 같다.

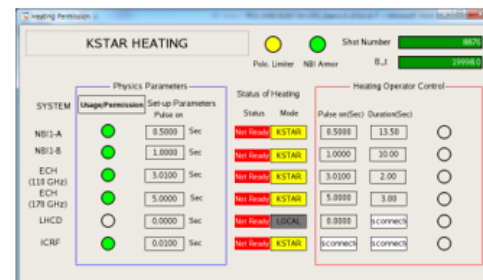


6. EDM 윈도우 오픈 (안 보일 경우)

- 개인 아이디 로그인에 아이디로 비밀번호를 입력한 것도 사용 가능
- OpIOS 서버 접속
- ssh -Y kstar@172.17.100.105 (kstarpcs 가 아니니 주의)
- yes -> password: k***2*1
- edm -x PCS-medm/PCS_status_display.edi

7. CPO 용 가열장치 Permission UI 오픈(PCS1)

- 아무 터미널이나 하나 열고 ssh 할 필요 없음
- "HeatingPermission CPO"라고 입력 (case sensitive???)하면 아래 창이 뜬다 (아래 그림은 CPO 용이 아니고 참조만 할 것)
- 모든 Heating permission "NO"로 설정



8. Webportal logbook 열기

- Logbook 오픈 -> <http://webportal.nfri.re.kr/logbook>

KSTAR campaign & PAC coordinator [2013-2014]

Rotation/Transport study shots

- KSTAR shots with CES/XICS/ECEI diagnostics collected

PUBLICATIONS ON EXPERIMENTS

- H.H. Lee et al., *Tearing modes induced by perpendicular electron cyclotron resonance heating in the KSTAR tokamak*, Nucl. Fusion 54(2014), 103008
- Jayhyun Kim et al., *Suppression of edge localized mode crashes by multi-spectral non-axisymmetric fields in the KSTAR*, Nucl. Fusion 57, 022001 (2017)
 - Type-I ELM crash suppression reproduced both consistent and inconsistent suppression performances when compared to the DIII-D results
 - indicates a dependency of ELM suppression on the heating level and the associated kink-like plasma responses
 - dominant and malign kink-like plasma responses over the benign gap filling effects

n=2 (mid) + n=1 (top/bot)

Port	L	P	D	H
ϕ	0°	90°	180°	270°
Top	+	-	-	+
Mid	+	-	+	-
Bot	+	+	-	-

➡ Proxy (n=1 NA field)
➡ Reference (n=2 NA field)
➡ Proxy (n=1 NA field)

$$\vec{B}_{\text{total}} = \vec{B}_{\text{eq}} + \vec{B}_{\text{int}} + \vec{B}_{\text{RMP}} + \vec{B}_{\text{pr}}$$

$$\vec{B}_{\text{pr}} = \vec{B}_{\text{pr}}^{\text{lin}} + \vec{B}_{\text{pr}}^{\text{nl}}$$

